Light Sources II

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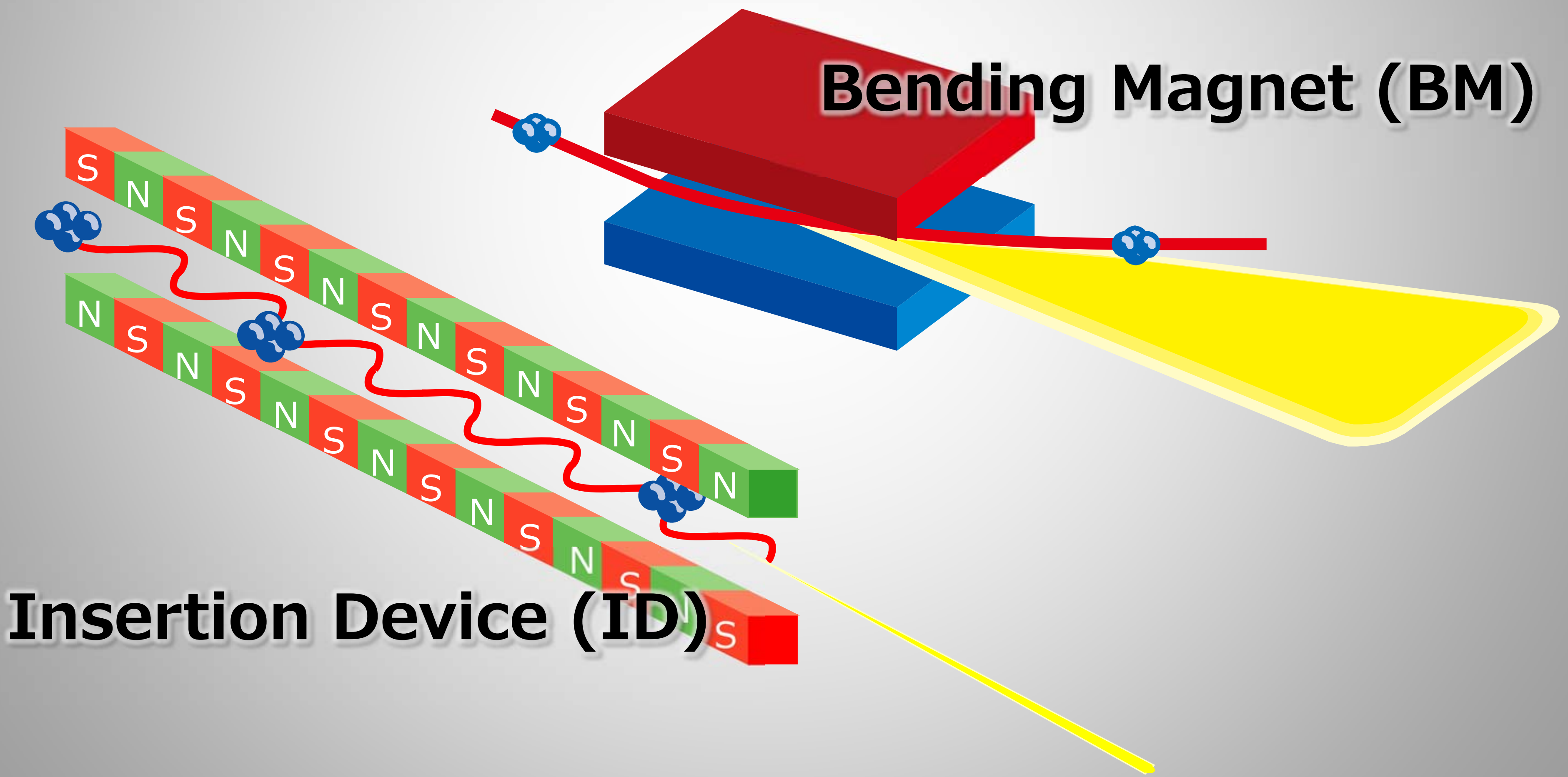
Outline

- Introduction
- Fundamentals of Light and SR
- **Overview of SR Light Source**
- Characteristics of SR (1)
- Characteristics of SR (2)
- Practical Knowledge on SR

What is SR Light Source?

Magnets to deflect the electron beam and generate SR.



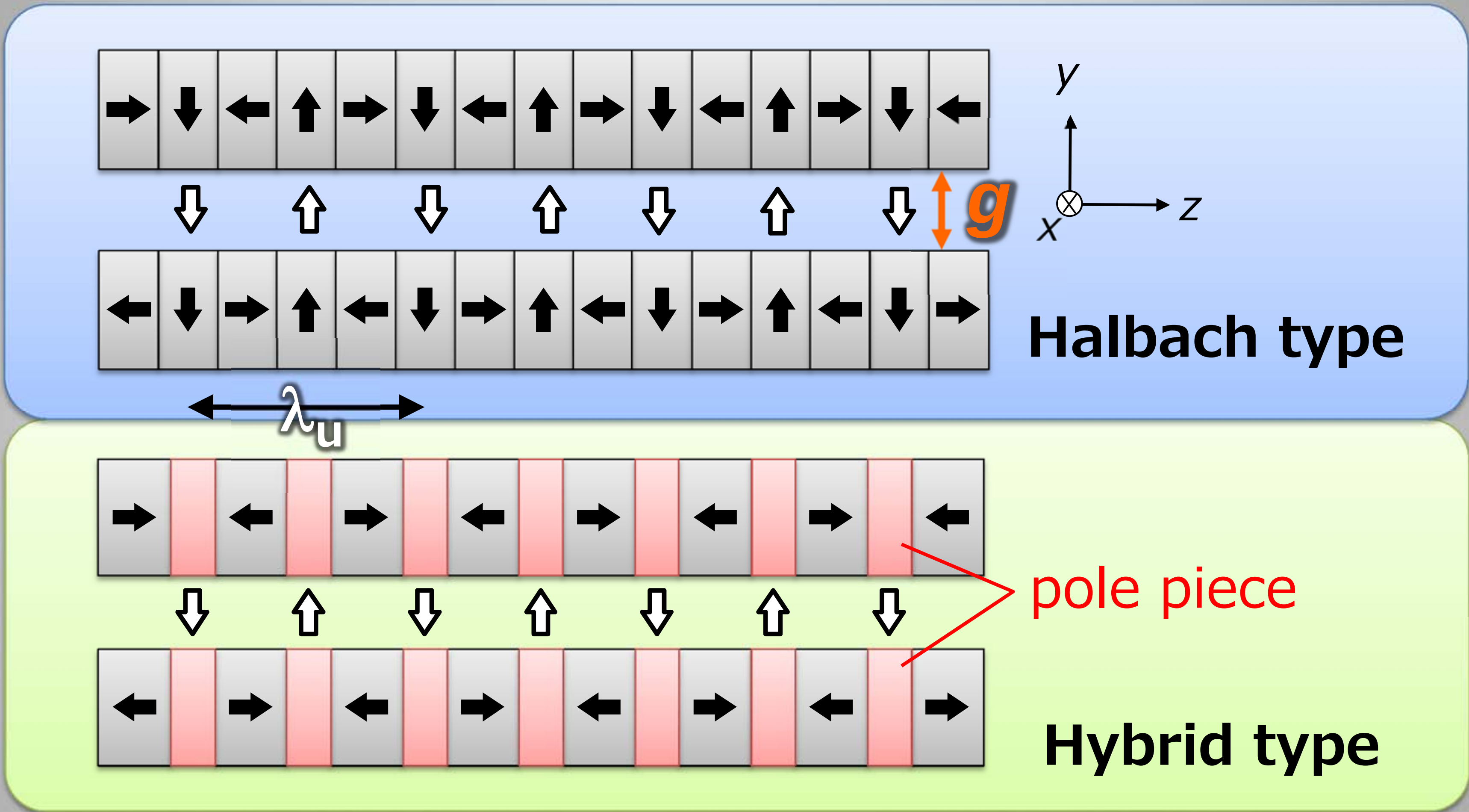
Bending Magnet

- One of the accelerator components in the storage ring.
- Generate **uniform field** to guide the electron beam into a **circular orbit**.
- EMs combined with highly-stable power supplies are adopted in most BMs to satisfy the stringent requirement on field quality and stability.
- Superconducting magnets are used in a few facilities in pursuit of harder x rays.

Insertion Device

- Installed (inserted) into the straight section of the storage ring between two adjacent BMs.
- Generate a **periodic magnetic field** to let the injected electron beam move along a **periodic trajectory**.
- Most IDs are composed of PMs, while EMs are used for special use such as helicity switching.
- Two types: **wiggler** and **undulator**

Magnetic Circuit of IDs

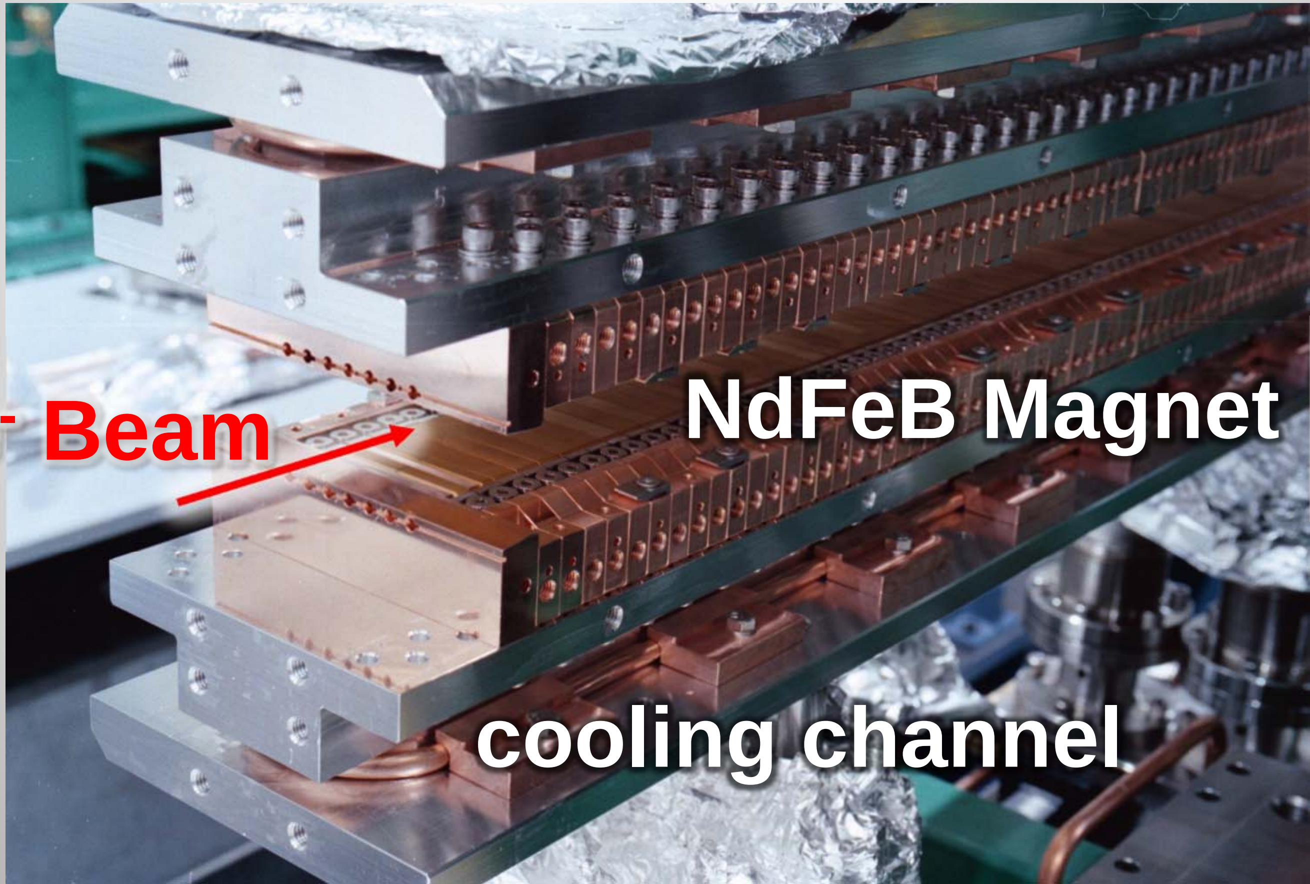


In each type, a sinusoidal magnetic field is obtained:

$$B_y(z) \sim B_0(B_r, g/\lambda_u) \sin\left(\frac{2\pi z}{\lambda_u}\right)$$

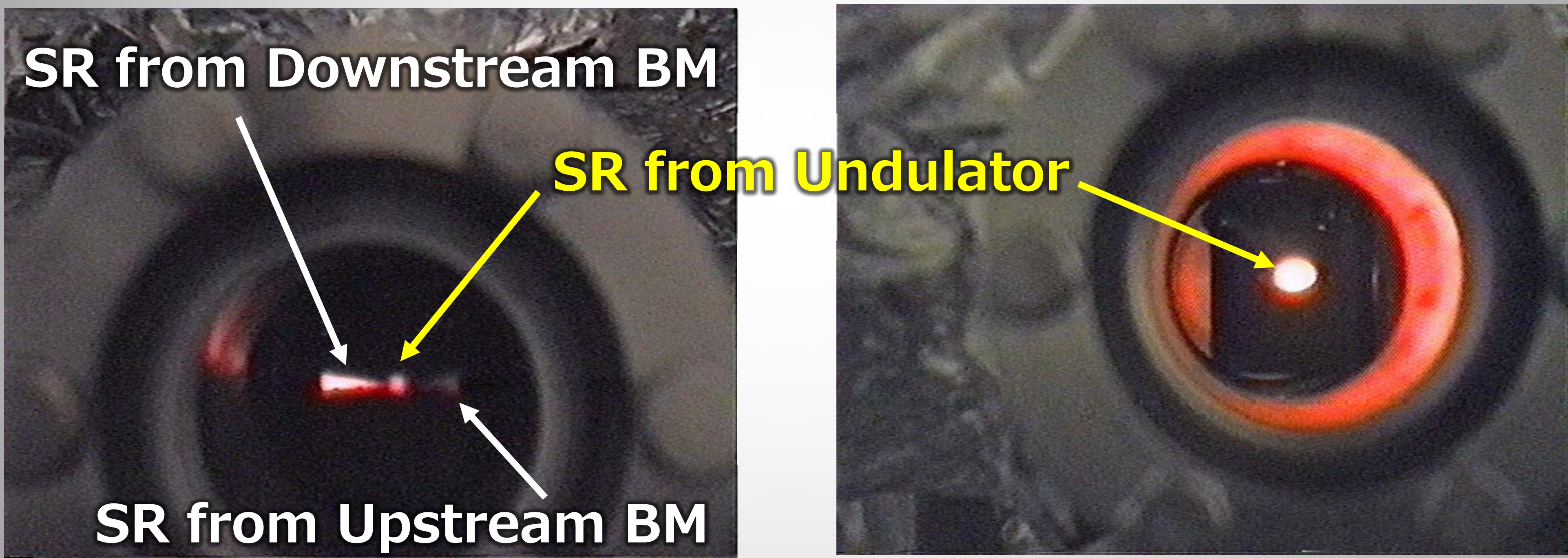
Example of ID Magnets

Halbach-type Magnet Array for
SPring-8 Standard Undulators



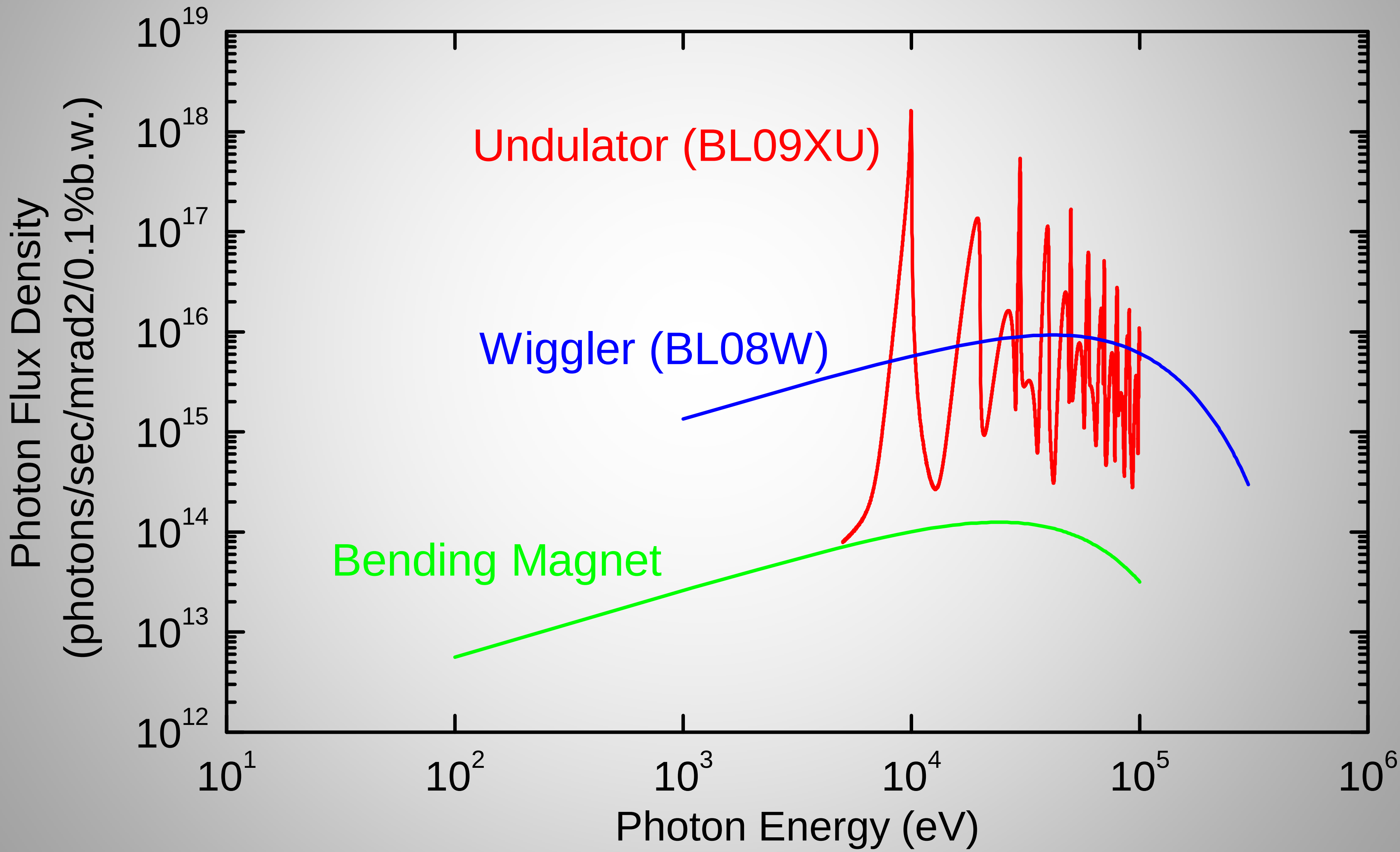
Example of SR Image

BL41XU@SP-8, first image of SR
with a fluorescent screen (<0.1mA)



***g* = 50 mm** **g = 20 mm**

Comparison of SR Light Sources



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  - **Radiation from Bending Magnets**
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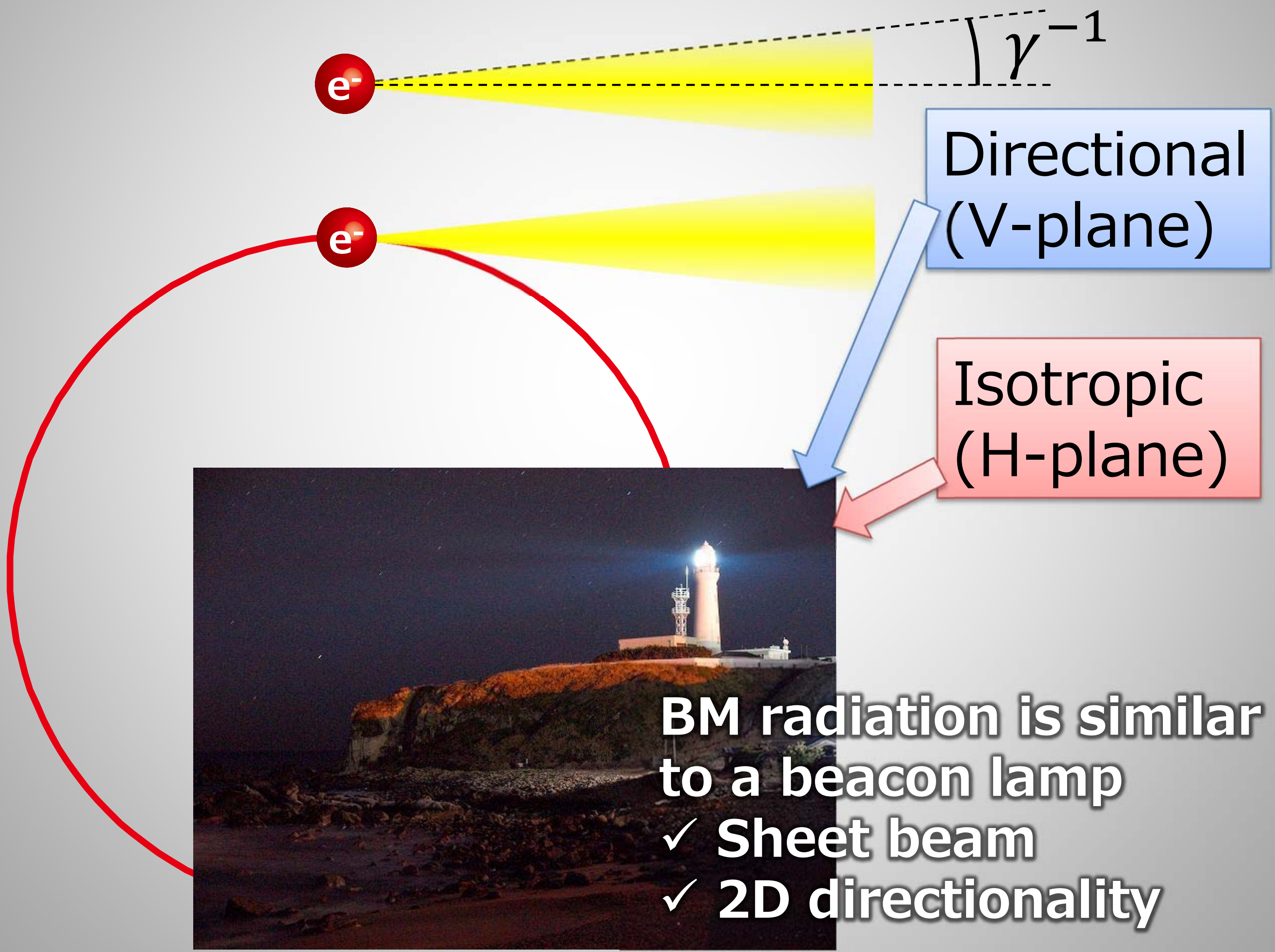
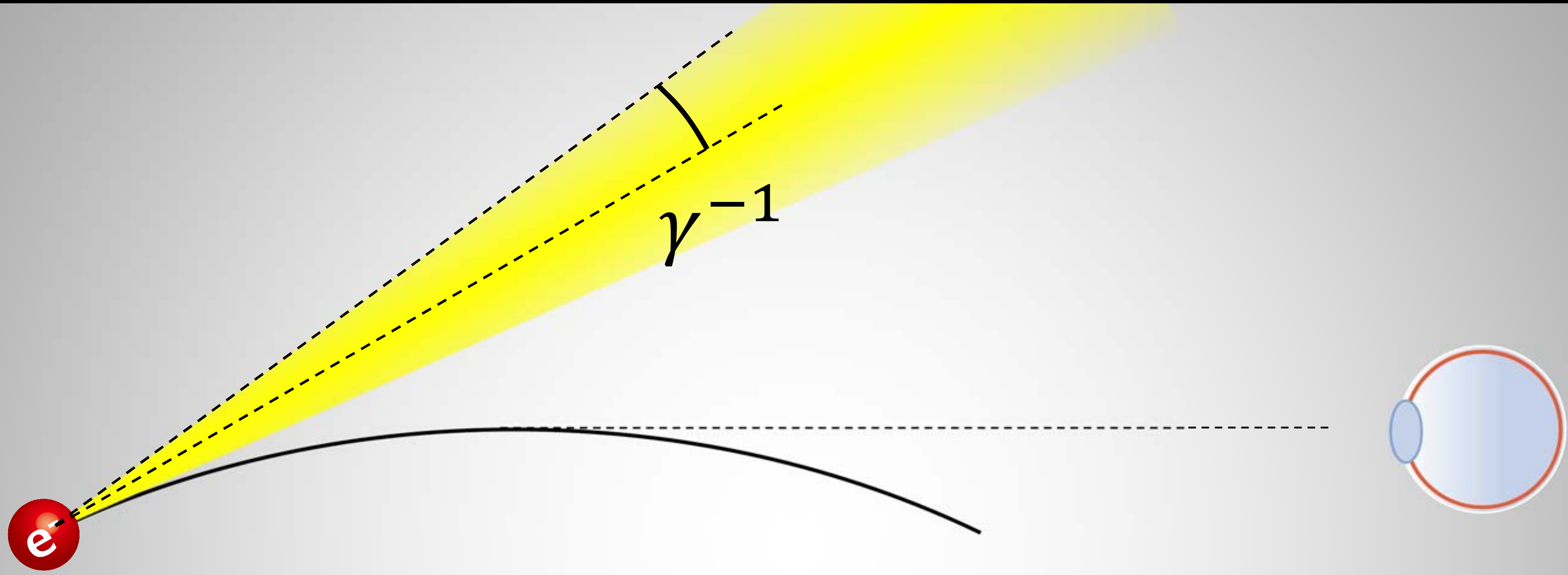
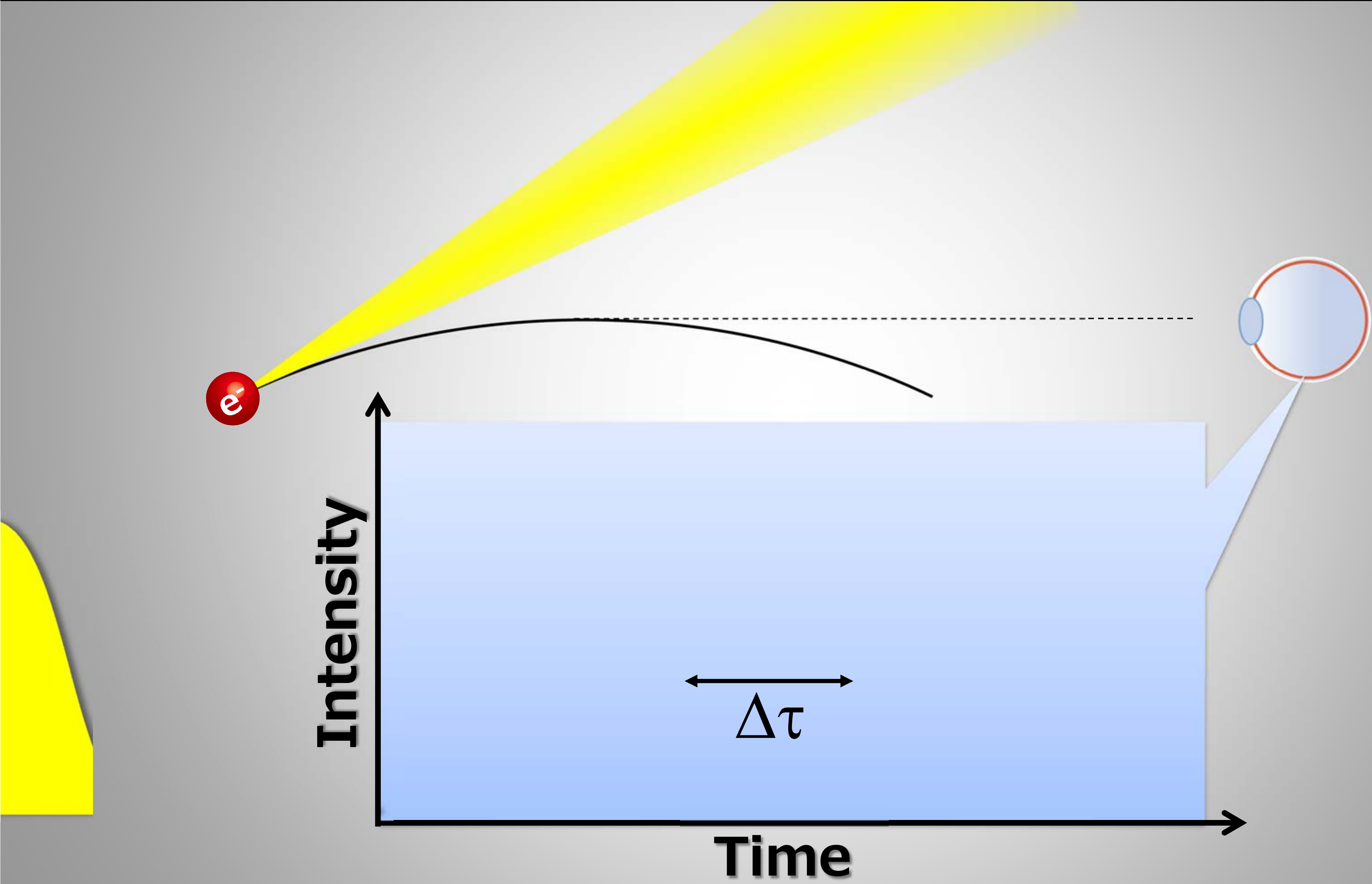
Directionality of BM Radiation



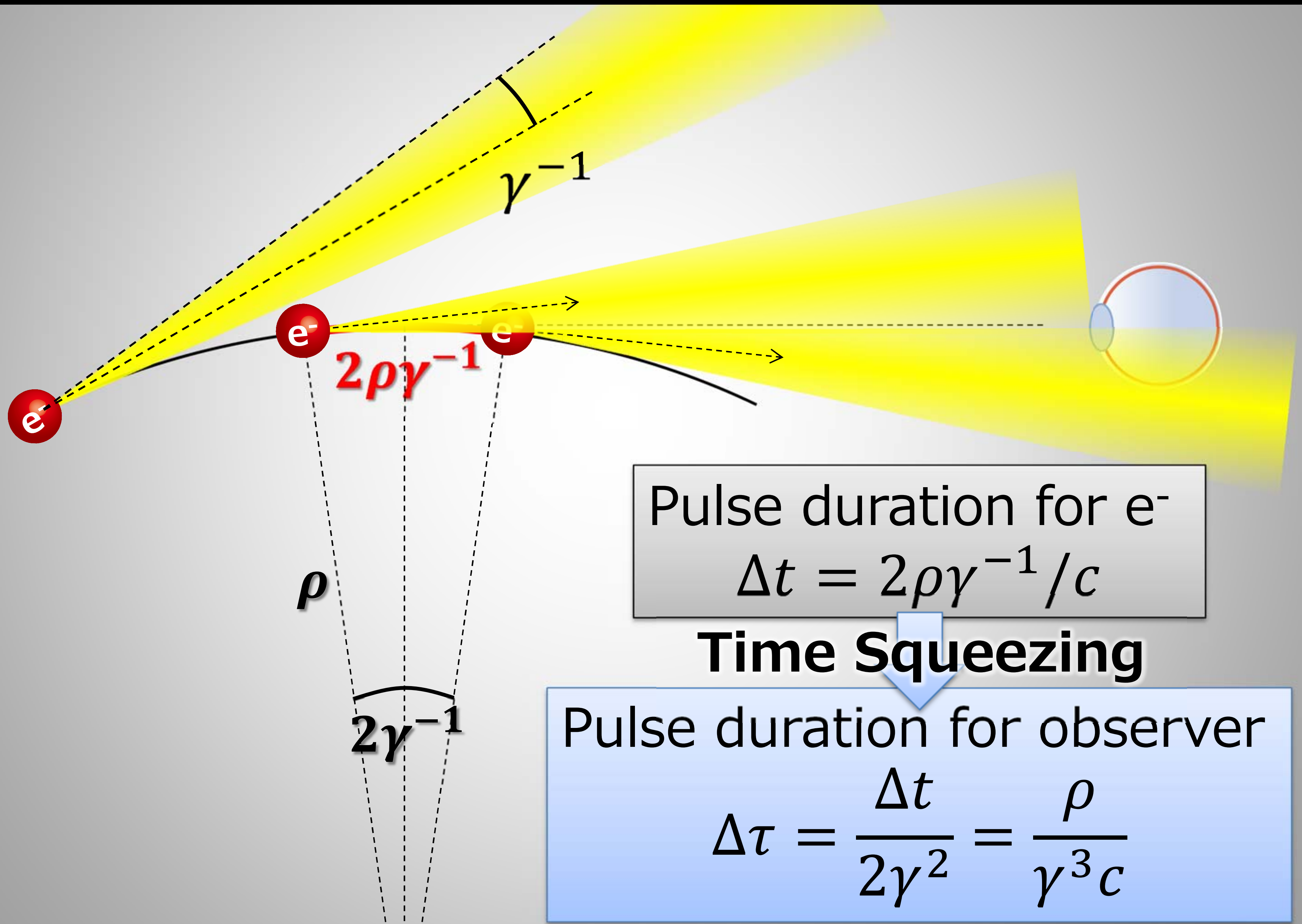
Pulse Structure of BM Radiation



Pulse Structure of BM Radiation



What's the Pulse Duration?



Pulse duration for e⁻

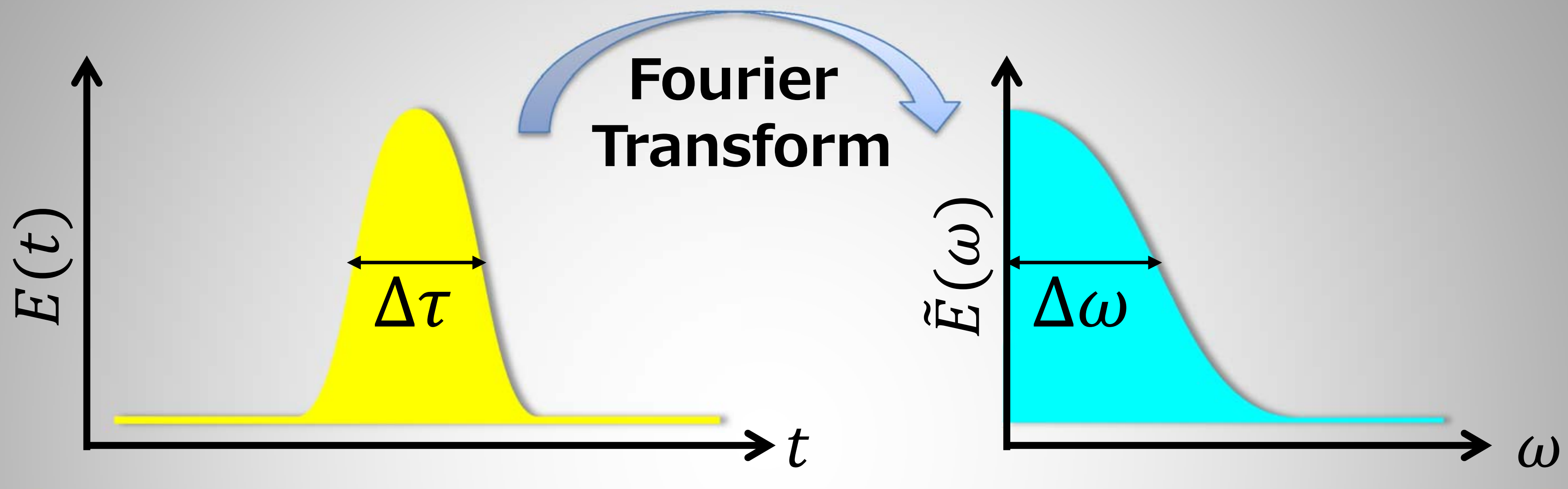
$$\Delta t = 2\rho\gamma^{-1}/c$$

Time Squeezing

Pulse duration for observer

$$\Delta\tau = \frac{\Delta t}{2\gamma^2} = \frac{\rho}{\gamma^3 c}$$

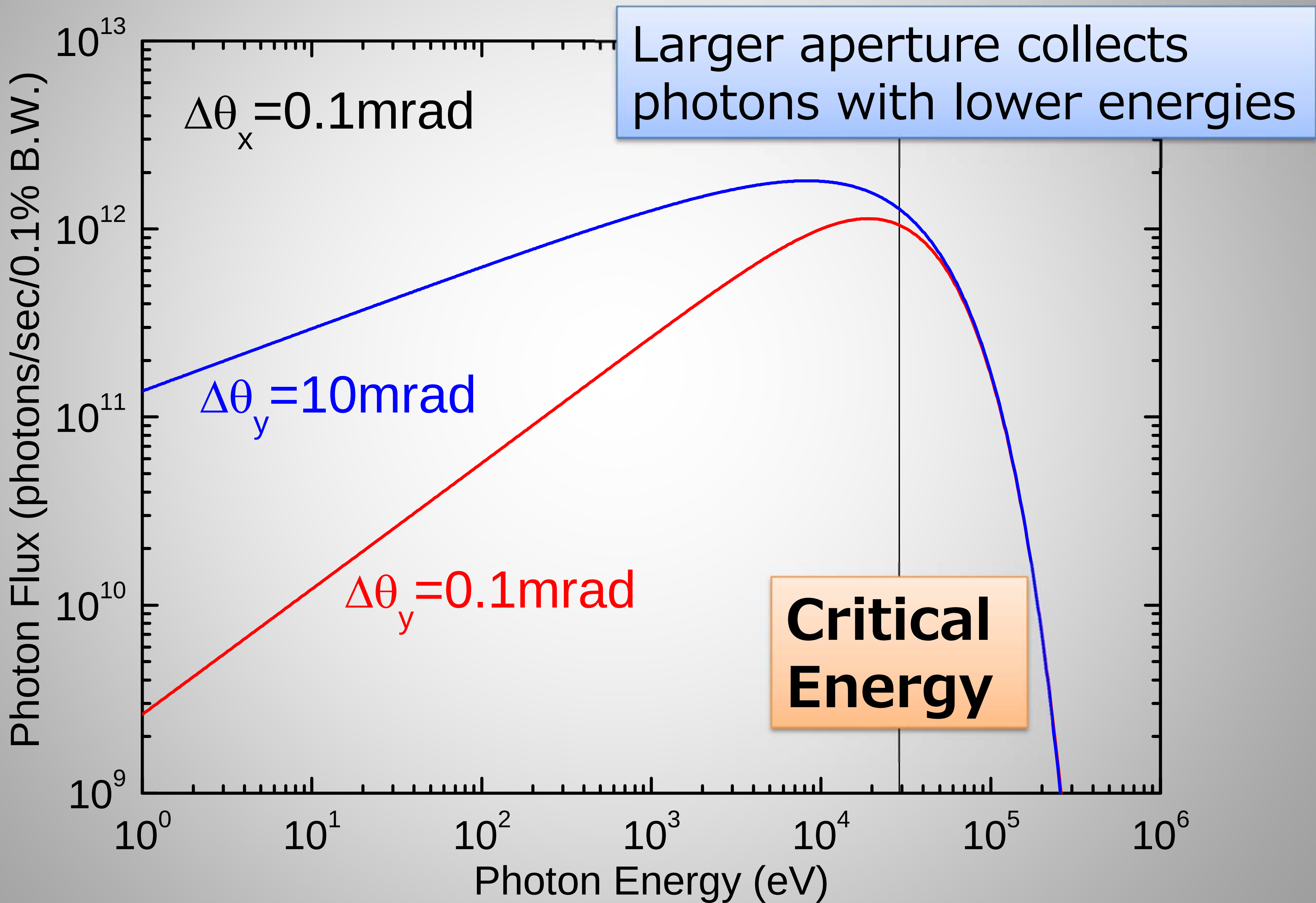
Spectrum of BM Radiation



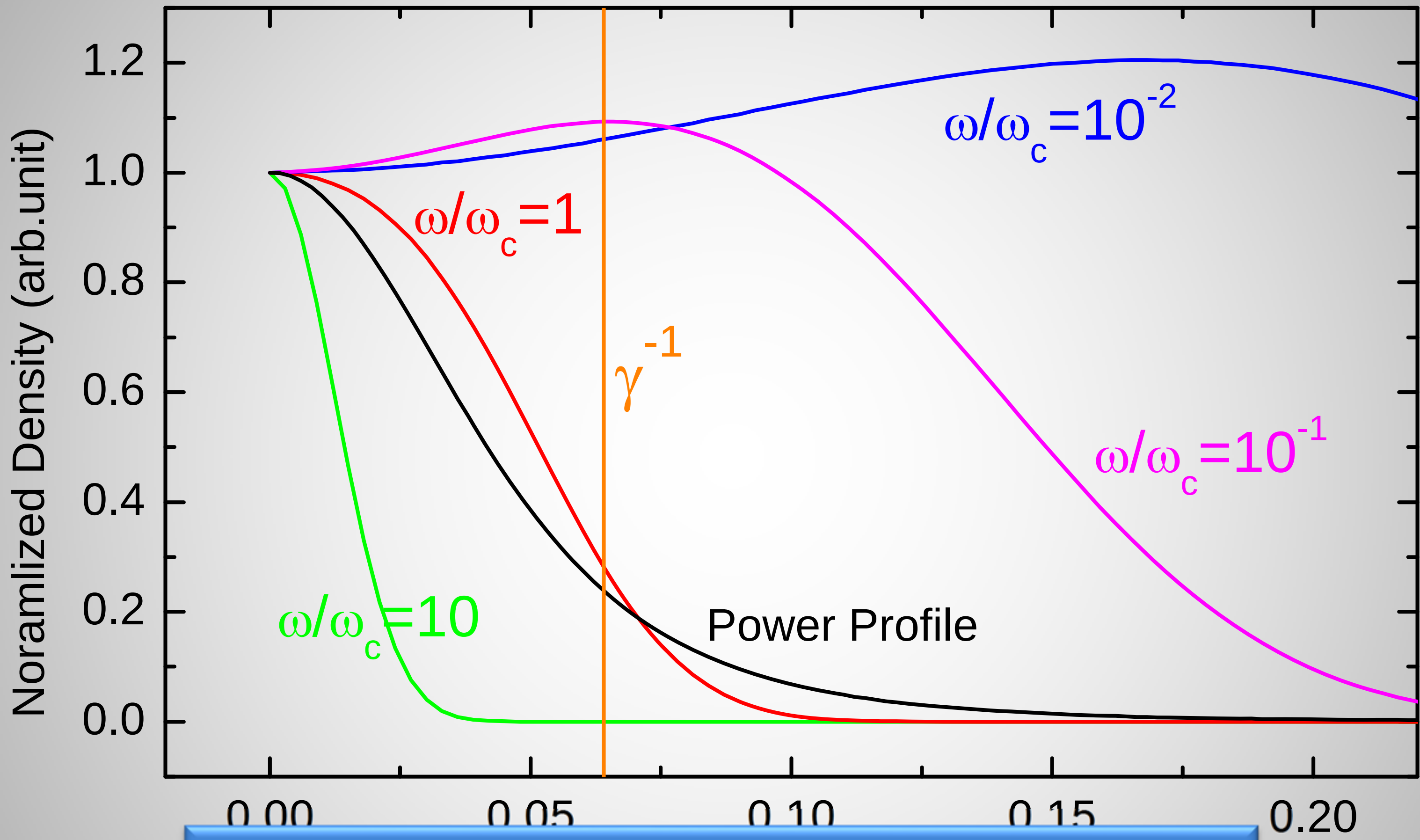
- ✓ BM radiation has a white spectrum reaching the frequency of $\Delta\omega \sim 1/\Delta\tau \sim \gamma^3 c/\rho$
- ✓ $\omega_c = \frac{3}{2\Delta\tau} = \frac{3\gamma^3 c}{2\rho}$ ("critical frequency") gives a criterion for BM spectrum.
- ✓ In practical units,

$$\hbar\omega_c(\mathrm{keV}) = 0.665 E^2\ (\mathrm{GeV}^2) B(\mathrm{T})$$

Example of Spectrum



Angular Profile of BM Radiation



Empirical formula for angular divergence of BM radiation

$$\sigma_{y'} = 0.6\gamma^{-1}\sqrt{\omega/\omega_c}$$

✓ pow

✓ large ... energy

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  - Undulator Radiation
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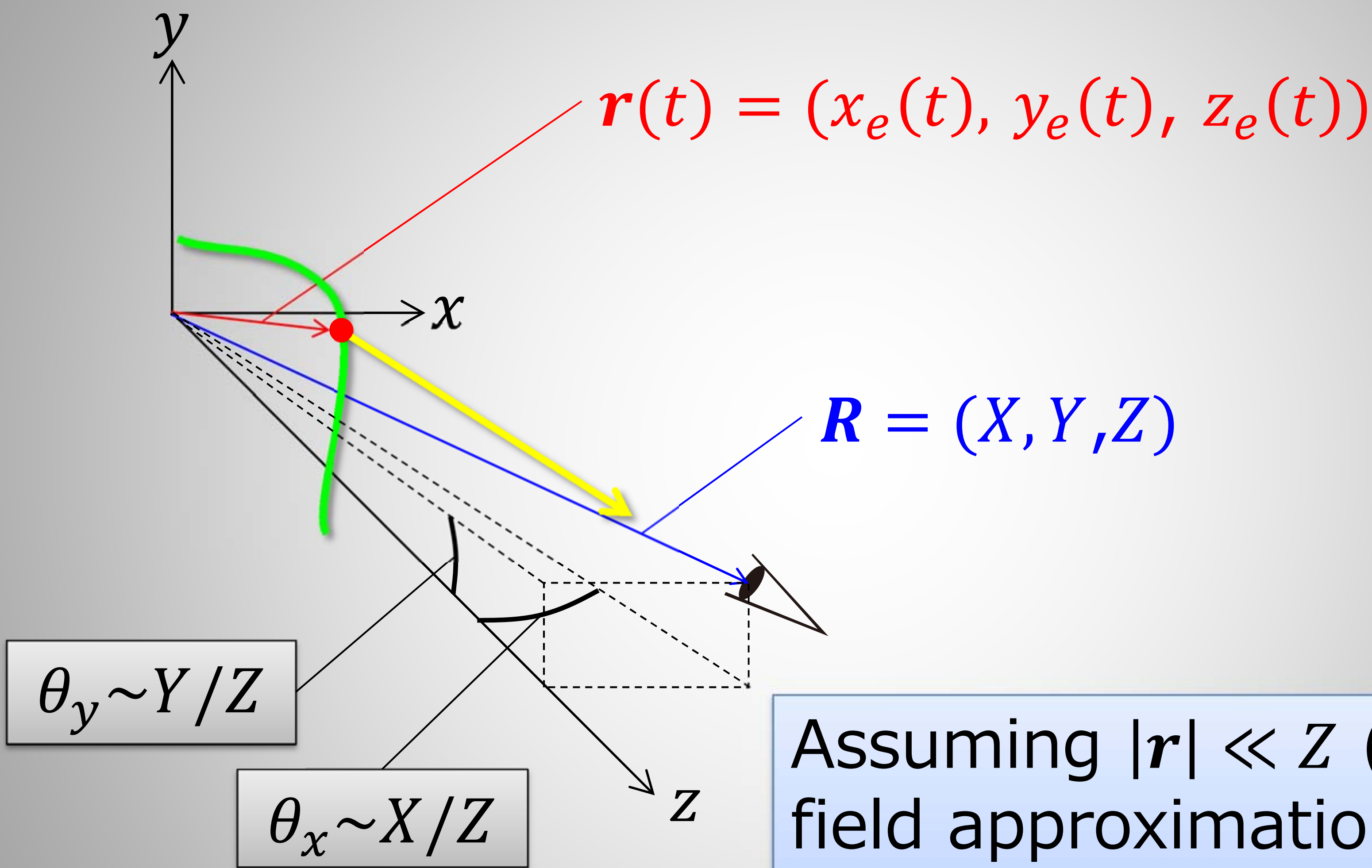
Coordinate Systems



Assuming $|\boldsymbol{r}| \ll Z$ (far-field approximation), the radiation pattern depends only on $\boldsymbol{\Theta} = (\theta_x, \theta_y)$

Field Integrals

$$\frac{d\boldsymbol{P}}{dt} = m\gamma\frac{d\boldsymbol{v}}{dt} = -e\boldsymbol{v}\times\boldsymbol{B} \Rightarrow \begin{cases} m\gamma\dot{v_x} = -e(v_yB_z - v_zB_y) \\ m\gamma\dot{v_y} = -e(v_zB_x - v_xB_z) \end{cases}$$

Equation of motion of an electron moving in a magnetic field ***B***

$$\Downarrow$$

$$m\gamma\frac{dv_{x,y}}{v_zdt} = m\gamma\frac{dv_{x,y}}{dz} = \pm eB_{y,x}$$

$$\beta_{x,y} = \pm\frac{e}{\gamma mc}\int^z B_{y,x}(z')dz' \equiv \pm\frac{e}{\gamma mc}I_{1y,1x}(z)$$

$$x_e, y_e = \pm\frac{e}{\gamma mc}\int^z dz'\int^{z'} B_{y,x}(z'')dz'' \equiv \pm\frac{e}{\gamma mc}I_{2y,2x}(z)$$

I_1, I_2: 1st and 2nd field integrals of the ID

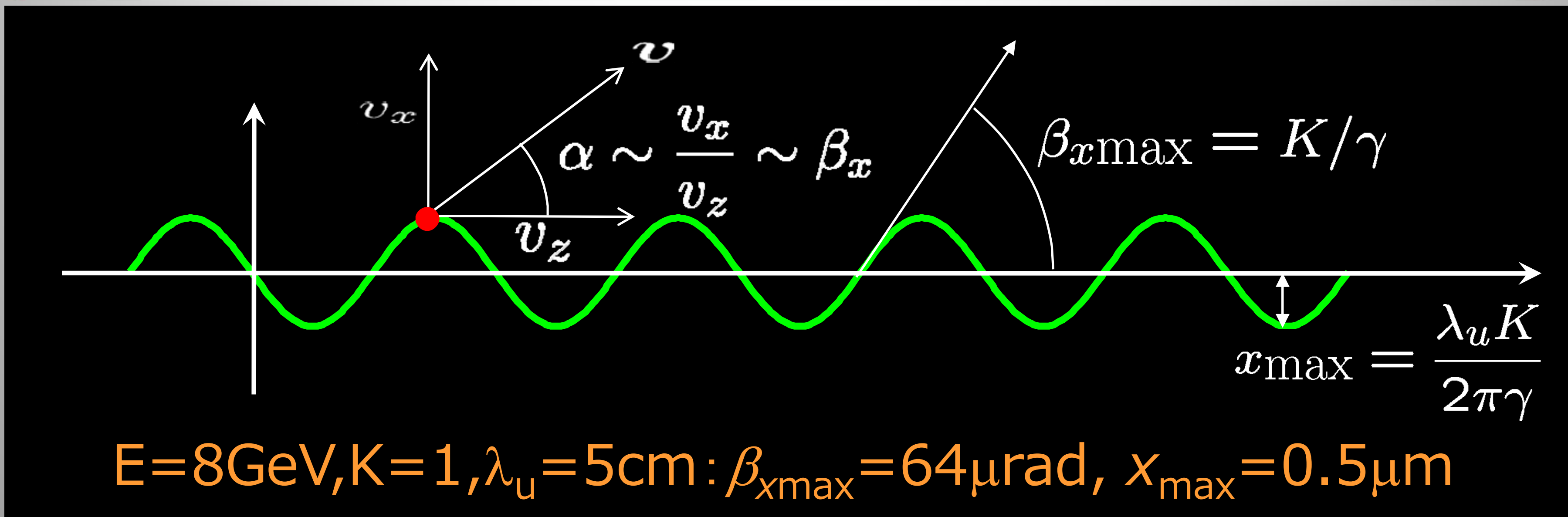
Trajectory in an Ideal ID

$$\begin{cases} B_x(z) = 0 \\ B_y(z) = B_0 \sin\left(\dfrac{2\pi z}{\lambda_u}\right) \end{cases} \quad \begin{cases} \beta_y(z) = 0 \\ \beta_x(z) = \dfrac{K}{\gamma}\cos\left(\dfrac{2\pi z}{\lambda_u}\right) \end{cases} \quad \begin{cases} y_e(z) = 0 \\ x_e(z) = \dfrac{\lambda_u K}{2\pi\gamma}\sin\left(\dfrac{2\pi z}{\lambda_u}\right) \end{cases}$$

Magnetic Field ⇒ **Velocity** ⇒ **Position**

$$K = \frac{eB_0\lambda_u}{2\pi mc} = 93.37 B_0(\mathrm{T})\lambda_u(\mathrm{m})$$

✓ **K Value**
✓ **Deflection Parameter**



Effects due to the ID Magnetic Field

transverse velocity

$$\beta_x(z) = \frac{K}{\gamma}\cos\left(\frac{2\pi z}{\lambda_u}\right)$$

longitudinal velocity

$$\beta_z = \sqrt{\beta^2 - \beta_x^2} = 1 - \frac{1}{2\gamma^2} - \frac{K^2}{4\gamma^2} - \frac{K^2}{4\gamma^2}\cos\left(\frac{4\pi z}{\lambda_u}\right)$$

Original value ($1 - \frac{1}{2\gamma^2}$)

Oscillating Term ($-\frac{K^2}{4\gamma^2}\cos\left(\frac{4\pi z}{\lambda_u}\right)$)

Average Value $\overline{\beta_z}$ ($1 - \frac{1}{2\gamma^2} - \frac{K^2}{4\gamma^2}$)

ID field induces:

- ✓ transverse (x) oscillation
- ✓ longitudinal (z) oscillation
- ✓ effective deceleration ($\Delta\beta_z = K^2/4\gamma^2$)

General Form of Time Squeezing

$$\frac{d\tau}{dt} = 1 - \boldsymbol{\beta} \cdot \boldsymbol{n}$$

$$\begin{aligned} \beta_z &= \sqrt{\beta^2 - \beta_x^2 - \beta_y^2} \\ &\sim 1 - (\gamma^{-2} + \beta_x^2 + \beta_y^2)/2 \\ n_z &\sim 1 - (\theta_x^2 + \theta_y^2)/2 \end{aligned}$$

$$= \frac{1}{2\gamma^2} + (\theta_x - \beta_x)^2 + (\theta_y - \beta_y)^2$$

Time squeezing takes place most significantly when the electron is moving in the direction of observation ($\boldsymbol{\beta} = \boldsymbol{\theta}$).

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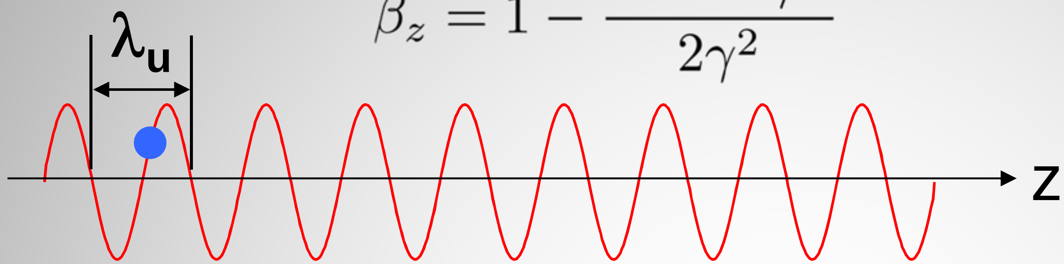
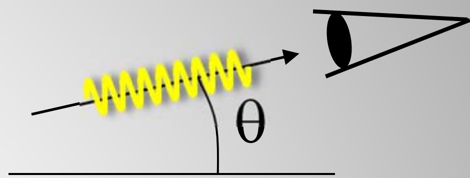
Fundamental Wavelength of UR

$$\bar{\beta}_z = 1 - \frac{1+K^2/2}{2\gamma^2}$$





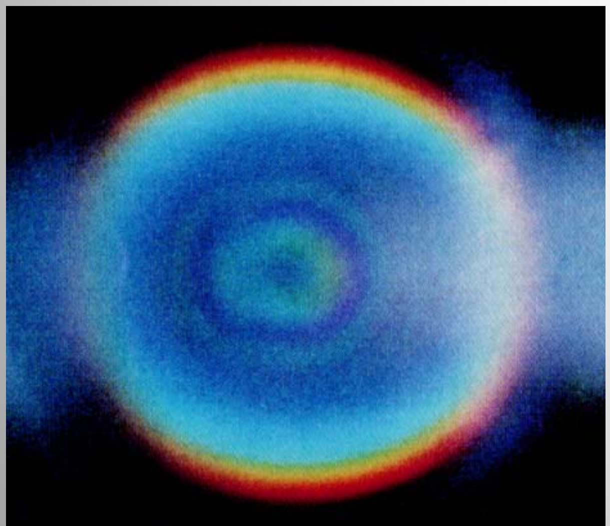
$$T = \lambda_u/v_z = \lambda_u/c$$

period of electron motion
= period of emitted light

Time Sequeezing

$$T' = T(1-\bar{\beta}_z\cos\theta)$$

period of observed light



H. Kitamura et al.,
J. Appl. Phys. 21 (1982) 1728

Fundamental Wavelength λ_1

$$\lambda_1 = \lambda_u(1-\bar{\beta}_z\cos\theta)$$

$$= \frac{\lambda_u}{2\gamma^2}(1+\gamma^2\theta^2+K^2/2)$$

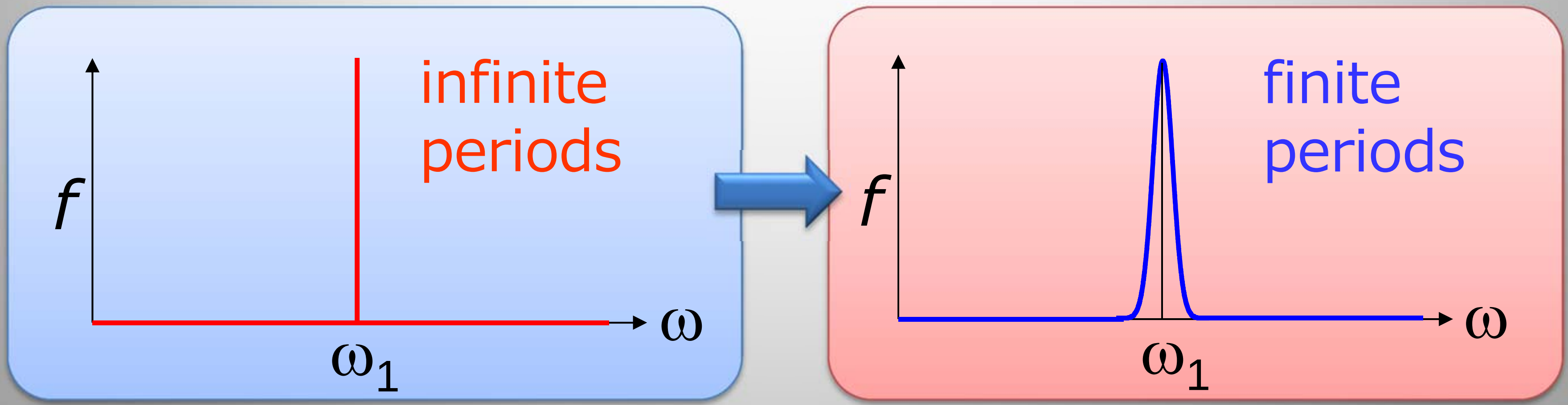
$$\omega_1 = 2\pi c/\lambda_1$$

UR with Infinite Periods

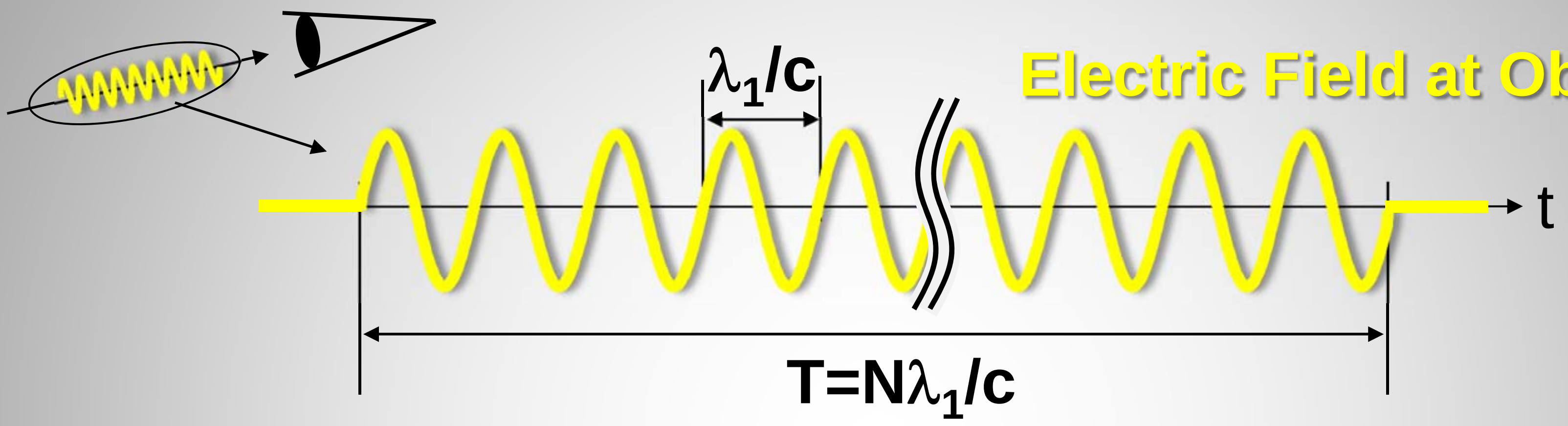
- If the undulator length is infinite, the pulse duration is infinitely long, and thus the radiation is completely monochromatic with line spectrum.

$$f(\theta, \omega) = \delta(\omega - \omega_1) = \delta\left(\omega - \frac{4\pi c\gamma^2/\lambda_u}{1 + K^2/2 + \gamma^2\theta^2}\right)$$

- In practice, the undulator length is finite, so the line spectrum is broadened.



Effects due to Finite Periods



$$E(t) = \begin{cases} E_0 \sin \omega_1 t & ; -T/2 \le t \le T/2 \\ 0 & ; t < -T/2,\ \ T/2 < t \end{cases} \quad , \quad \omega_1 = 2\pi c/\lambda_1$$

Fourier Transform

$$f(\theta,\omega) \propto |\tilde{E}(\omega)|^2 \propto \mathrm{sinc}^2 \left[\pi N \frac{\omega - \omega_1(\theta)}{\omega_1(\theta)} \right]$$

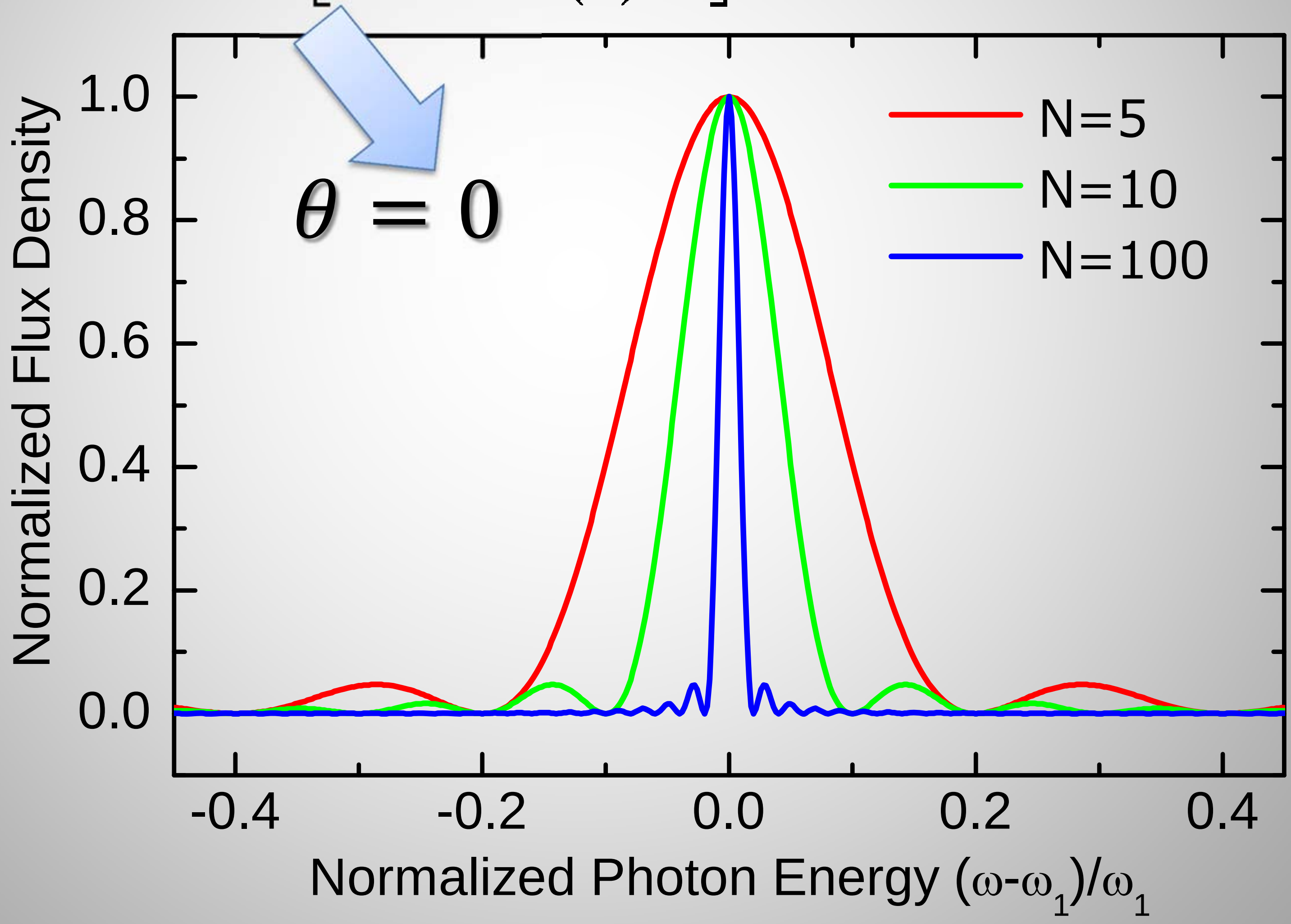
Square of "sinc" function dominates UR

Brief Note on UR Formulae

- In the previous derivations of UR spectral function, no knowledge on electrodynamics is required.
- In practice, E_0 is a complicated function of θ and K, and needs to be calculated by Fourier transforming the electric field derived from the Lienard-Wiecherd potential.
- However, the simple derivation gives us a clear understanding on UR properties.

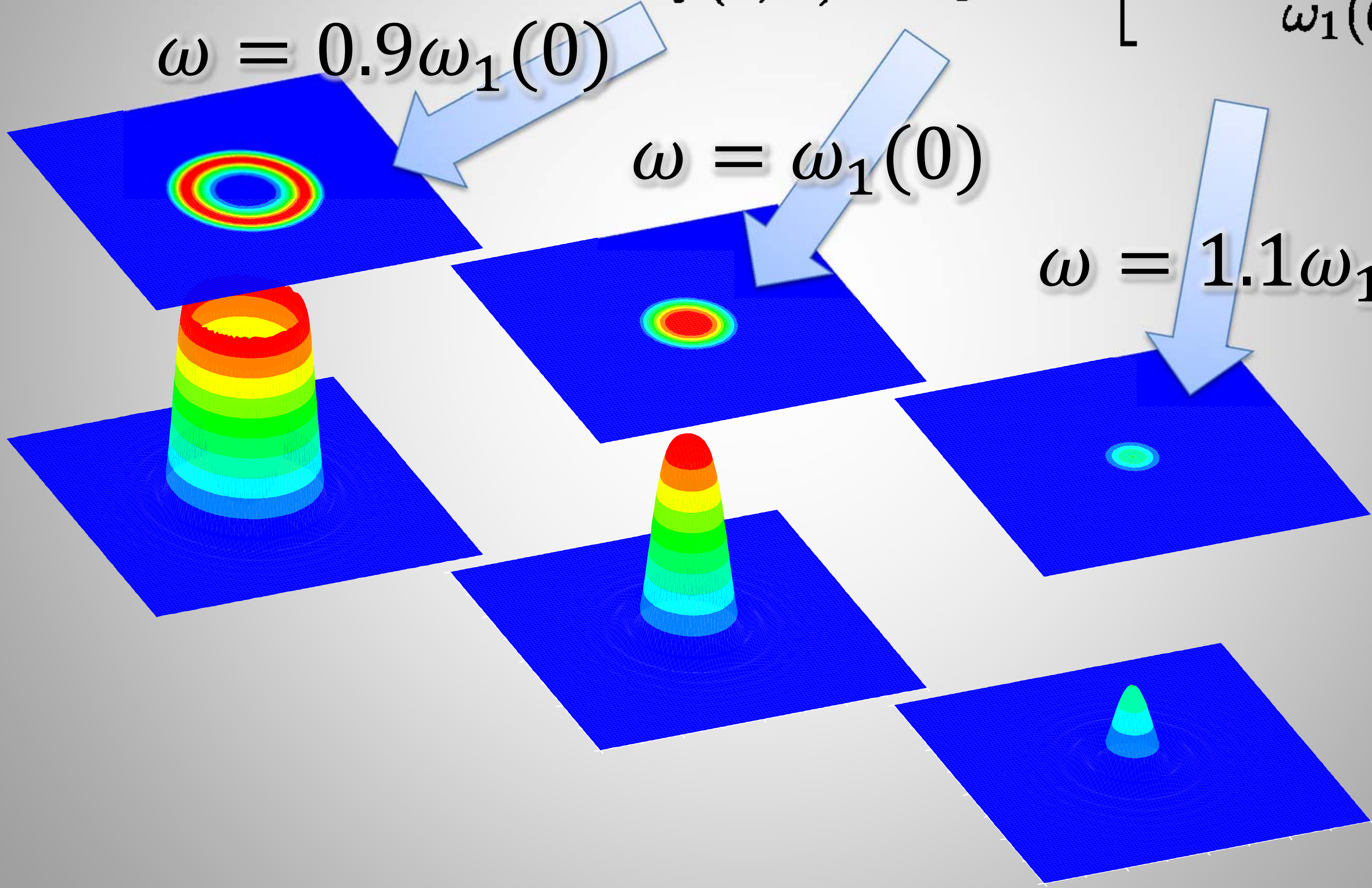
Energy Spectrum of UR

$$f(\theta,\omega) = F_0 \mathrm{sinc}^2\left[\pi N \frac{\omega - \omega_1(\theta)}{\omega_1(\theta)}\right]$$



Angular Profile of UR

$$f(\theta,\omega) = F_0 \mathrm{sinc}^2 \left[\pi N \frac{\omega - \omega_1(\theta)}{\omega_1(\theta)} \right]$$

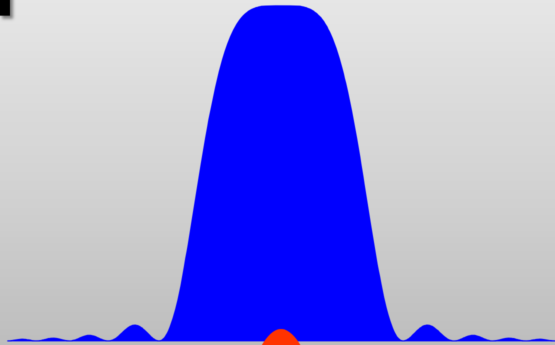


Angular Divergence of UR

UR is not a Gaussian Beam

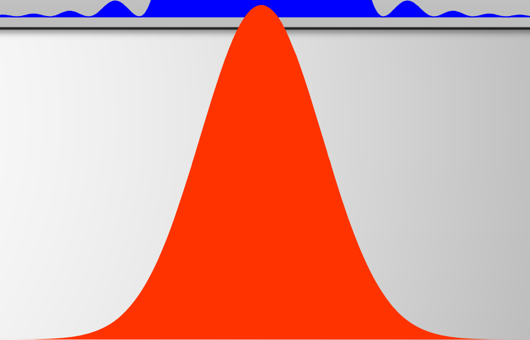
Angular Profile at $\omega=\omega_1(0)$

$$f(\theta,\omega) = F_0 \mathrm{sinc}^2\left[\frac{\pi N(\gamma\theta)^2}{1+K^2/2}\right]$$



Gaussian Profile with $\sigma_{r'}$

$$f_G(\theta) = F_0 \exp\left(-\frac{\theta}{2\sigma_r'^2}\right)$$



$$\int f(\theta,\omega)d\boldsymbol{\theta} = \int f_G(\theta)d\boldsymbol{\theta}$$

$$\sigma_{r'} = \sqrt{\frac{1+K^2/2}{4N\gamma^2}} = \sqrt{\frac{\lambda_1}{2L}}$$

"Natural" Angular Divergence of UR ($L = N\lambda_u$)

Source Size of UR

$$\sigma_{r'} = \sqrt{\frac{1 + K^2/2}{4N\gamma^2}} = \sqrt{\frac{\lambda_1}{2L}}$$

"Natural" Angular Divergence of UR ($L = N\lambda_u$)

✓ Unlike WR, UR is spatially coherent, i.e., diffraction limited.

✓ This means that $\sigma_r \sigma_{r'} = \lambda/4\pi$

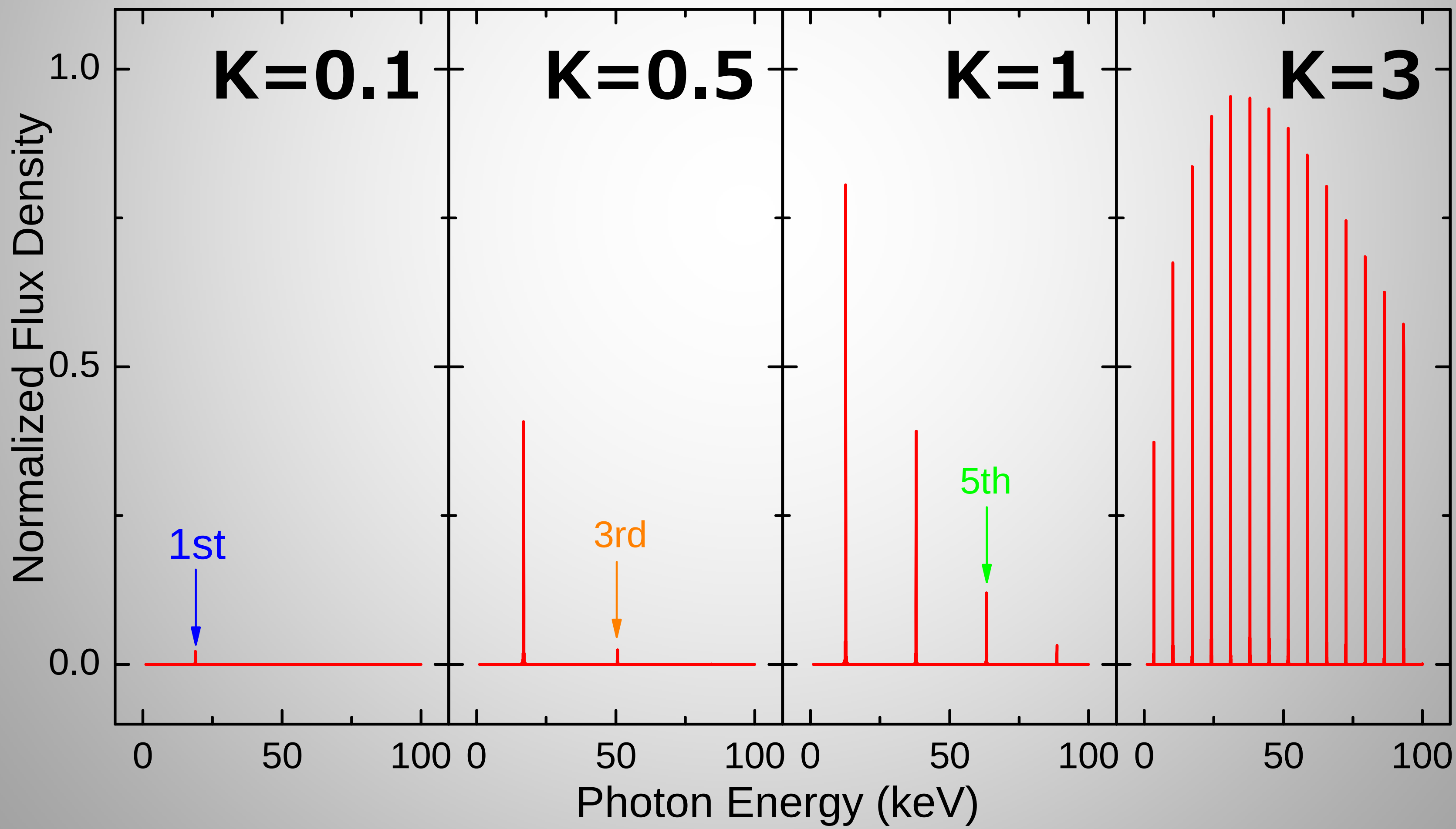
$$\sigma_r = \frac{\lambda_1}{4\pi\sigma_{r'}} = \frac{\sqrt{2L\lambda_1}}{4\pi}$$

"Natural" Source size of UR

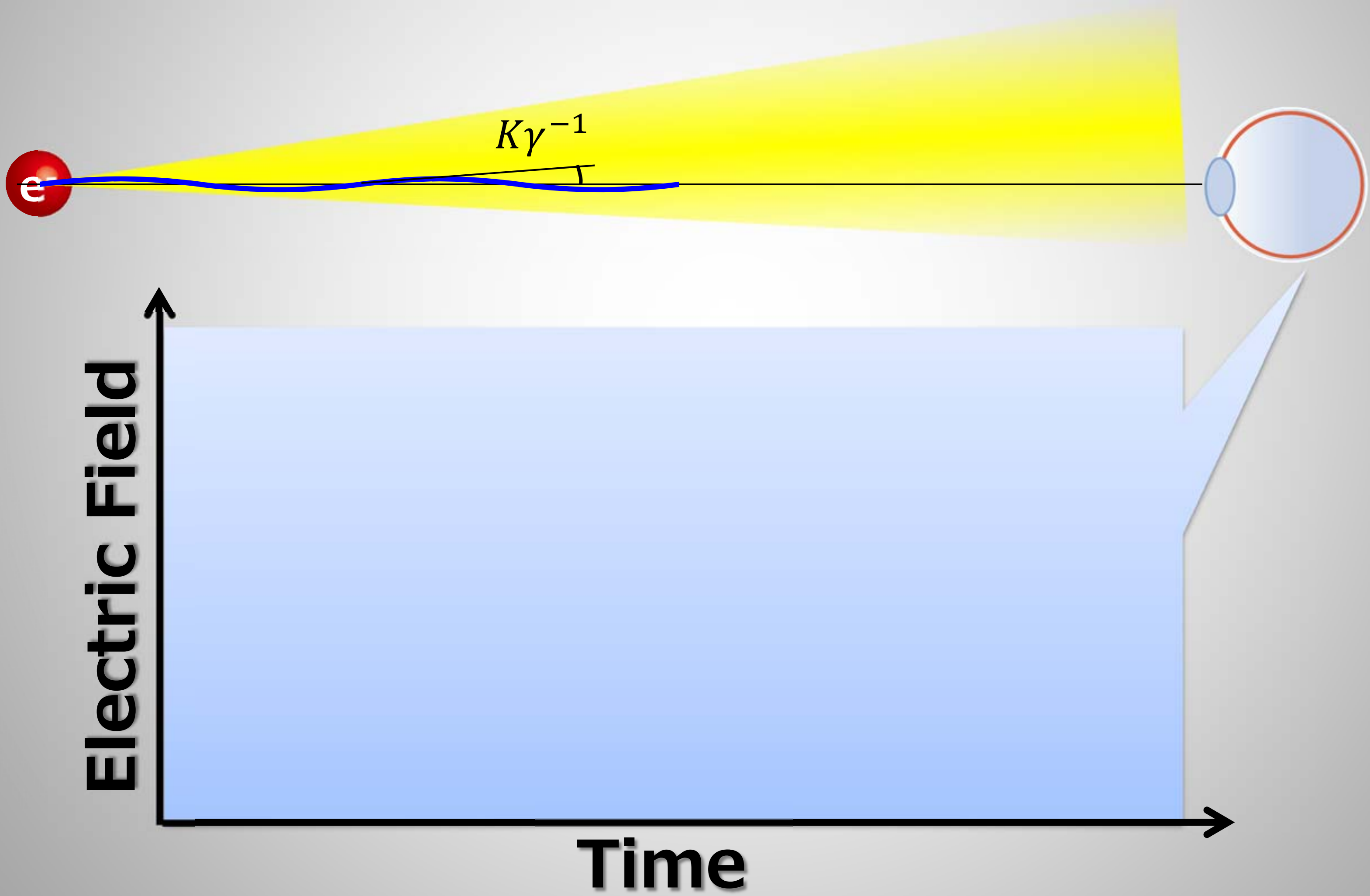
Longer device results in smaller angular divergence & larger source size, but the emittance does not change.

Higher Harmonics

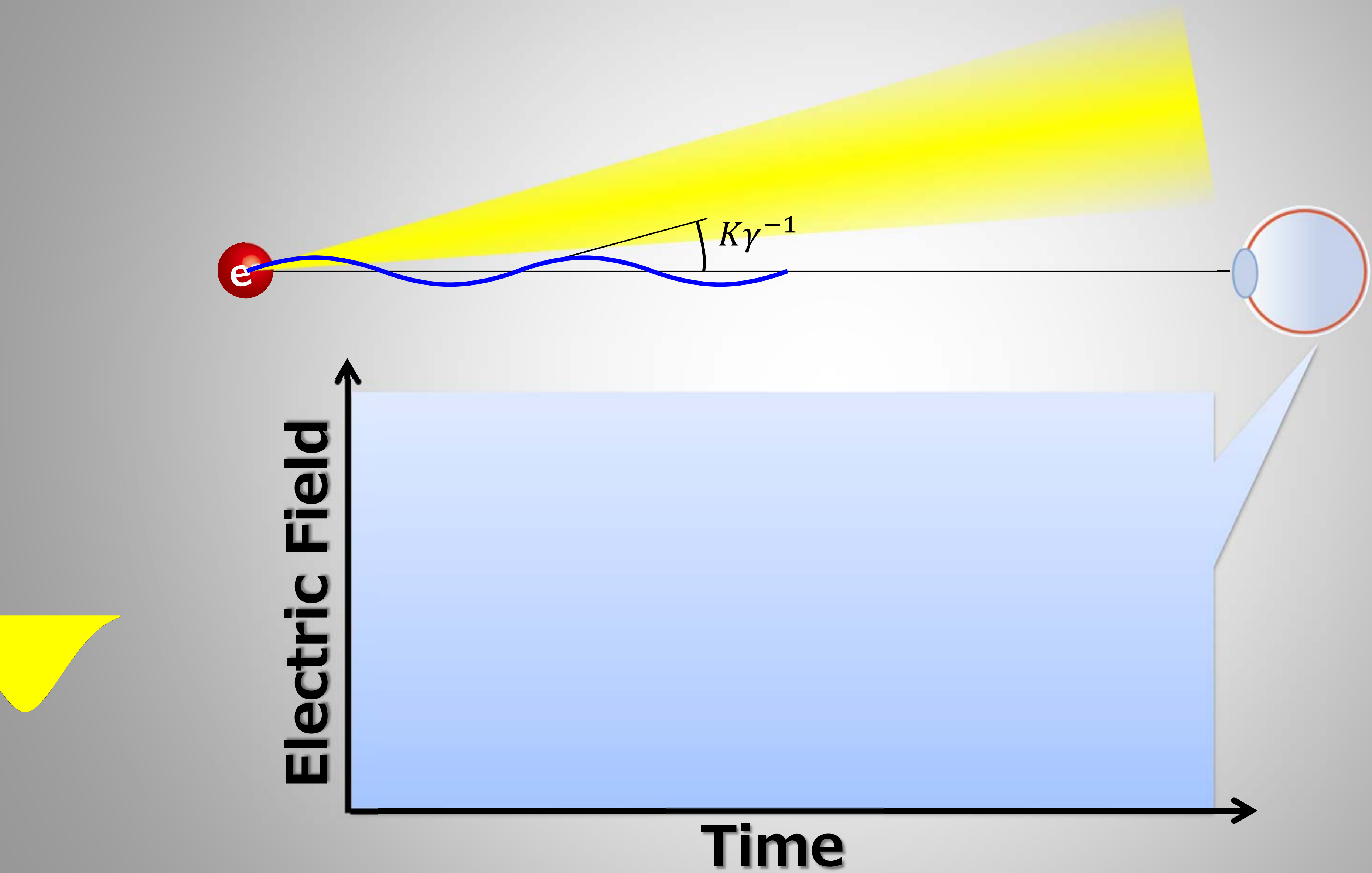
Photons with at $n\omega_1$ are observed as well as at ω_1, where n is an integer.



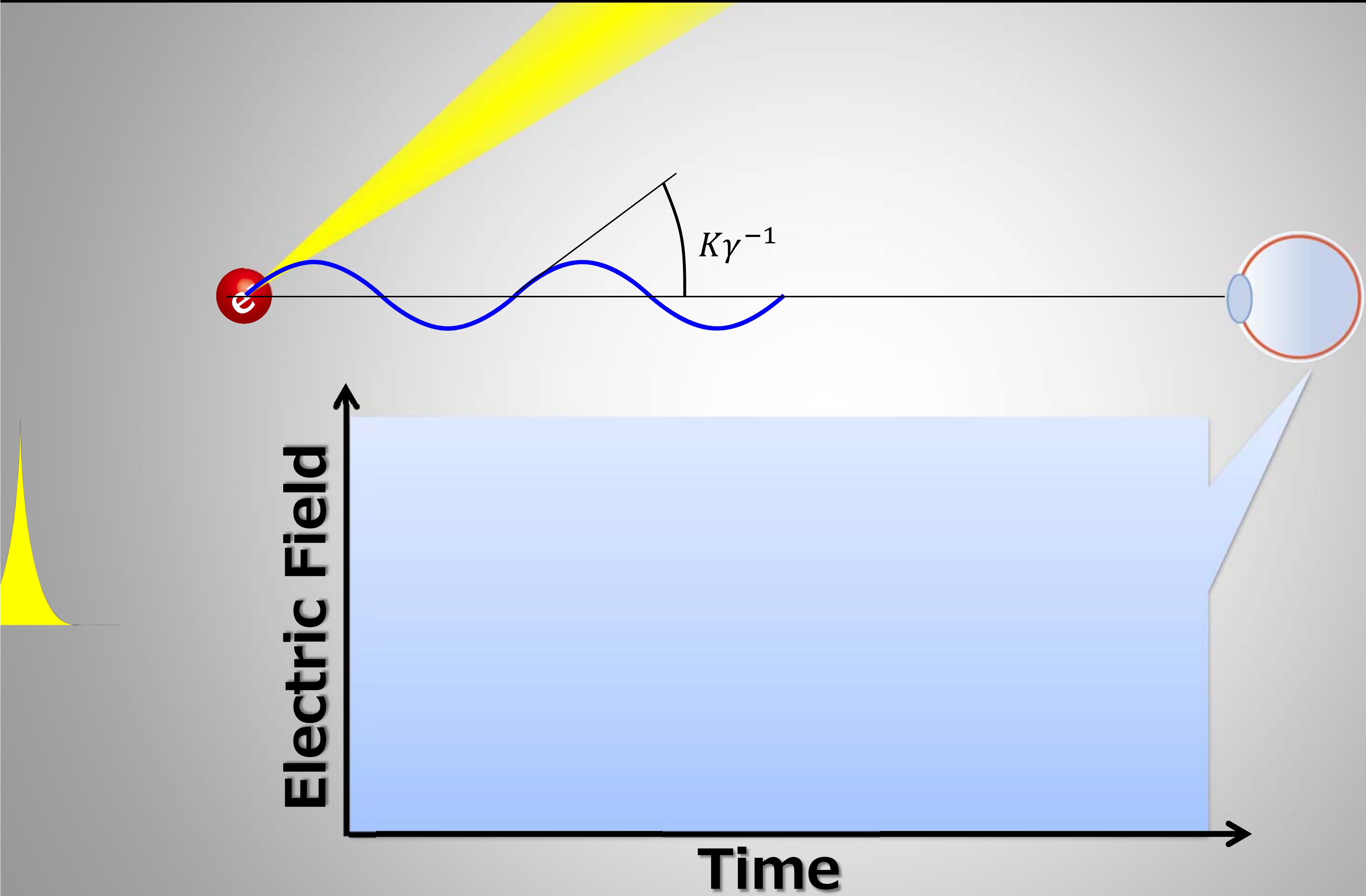
K << 1 Case



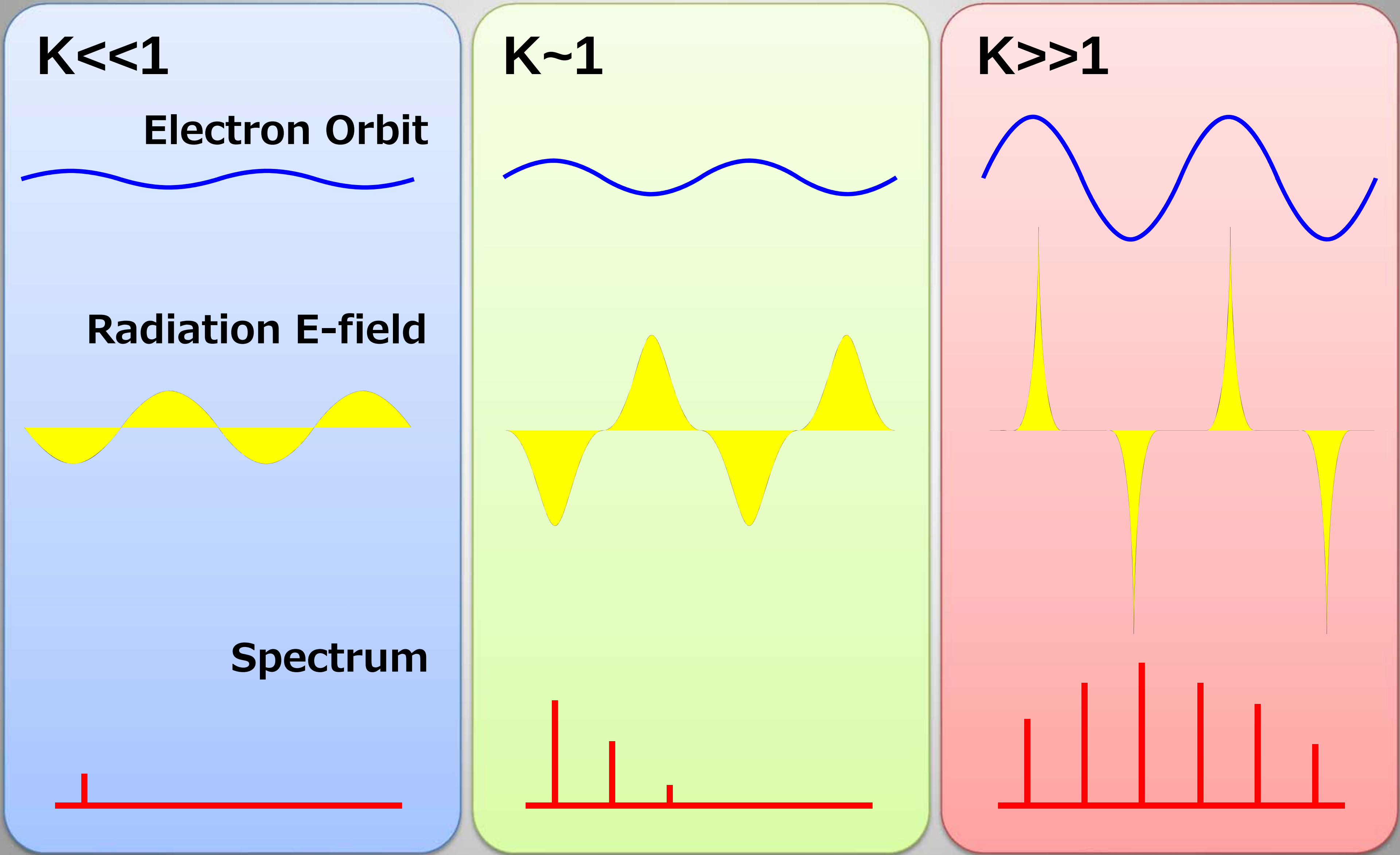
K ~ 1 Case



K >> 1 Case



Mechanisms of Higher Harmonics



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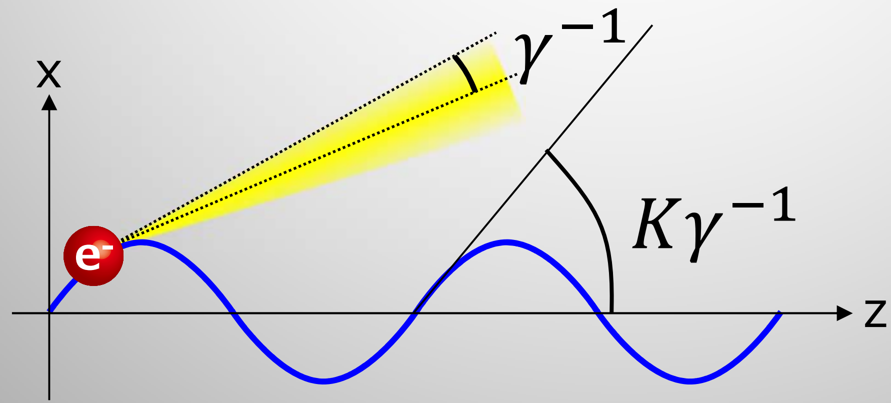
Wiggler Radiation

- In terms of the magnetic structure, wigglers are identical to undulators.
- Unlike UR, wiggler radiation (WR) is regarded as **incoherent sum of SR** emitted at each magnet pole.
  - Summation as photons
- As a result, WR is similar to BM radiation in many points, but some differences are present.

Comparison with BM

Item	BM	Wiggler
Spectral Profile	White (mooth)	Quasi-White (TBDL)
Vertical Directionality	Scales as γ^{-1} $\sigma_{y'} \sim 0.6\gamma^{-1}\sqrt{\omega/\omega_c}$	
Horizontal Directionality	Isotropic	$\sigma_{x'} \sim \frac{\sqrt{1+K^2/2}}{\gamma}$



SR cone angle (γ^{-1}) & deflection angle ($K\gamma^{-1}$)

Why Not Undulators?

- In terms of photon flux, coherence, heat load, etc., undulators are much superior to wigglers in most cases.
- However, wigglers are installed in specific beamlines, in order to obtain:
  – much higher photon energies than what are available with undulators
  – white spectrum for dispersive experiments
- “Superconducting wigglers” are sometimes used to enhance the field strength (~10T).

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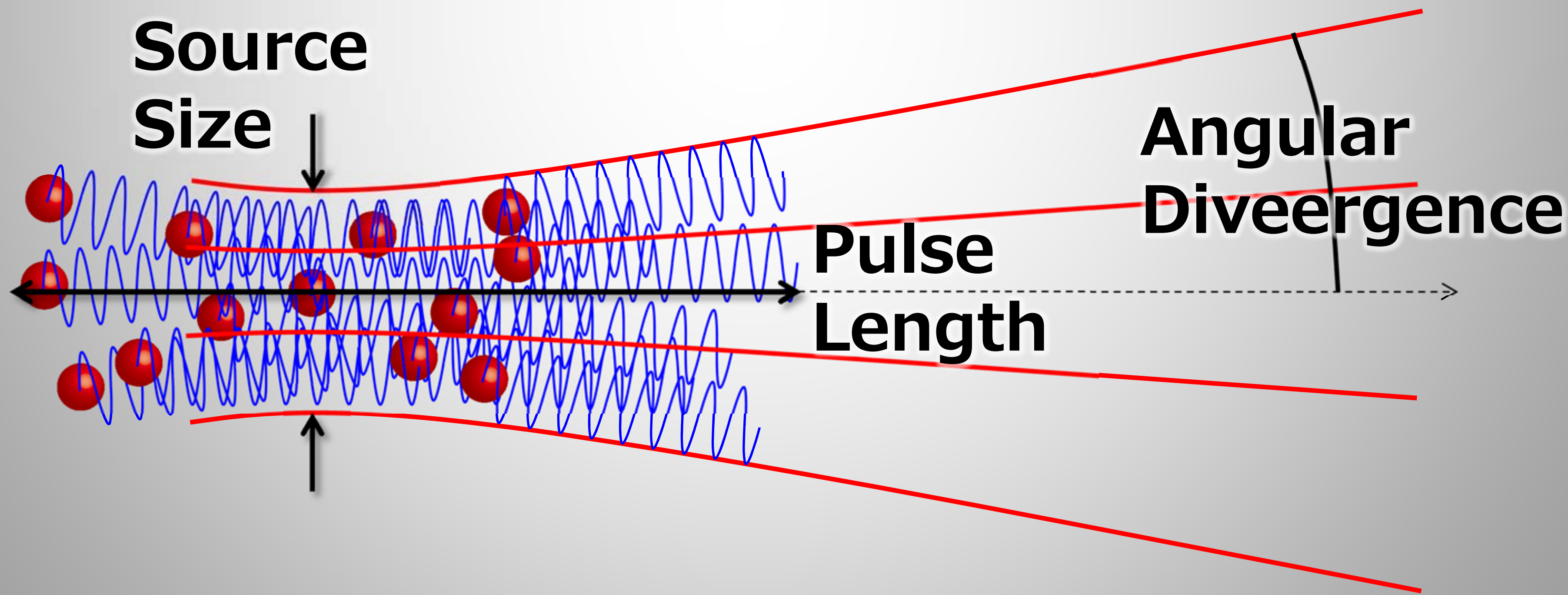
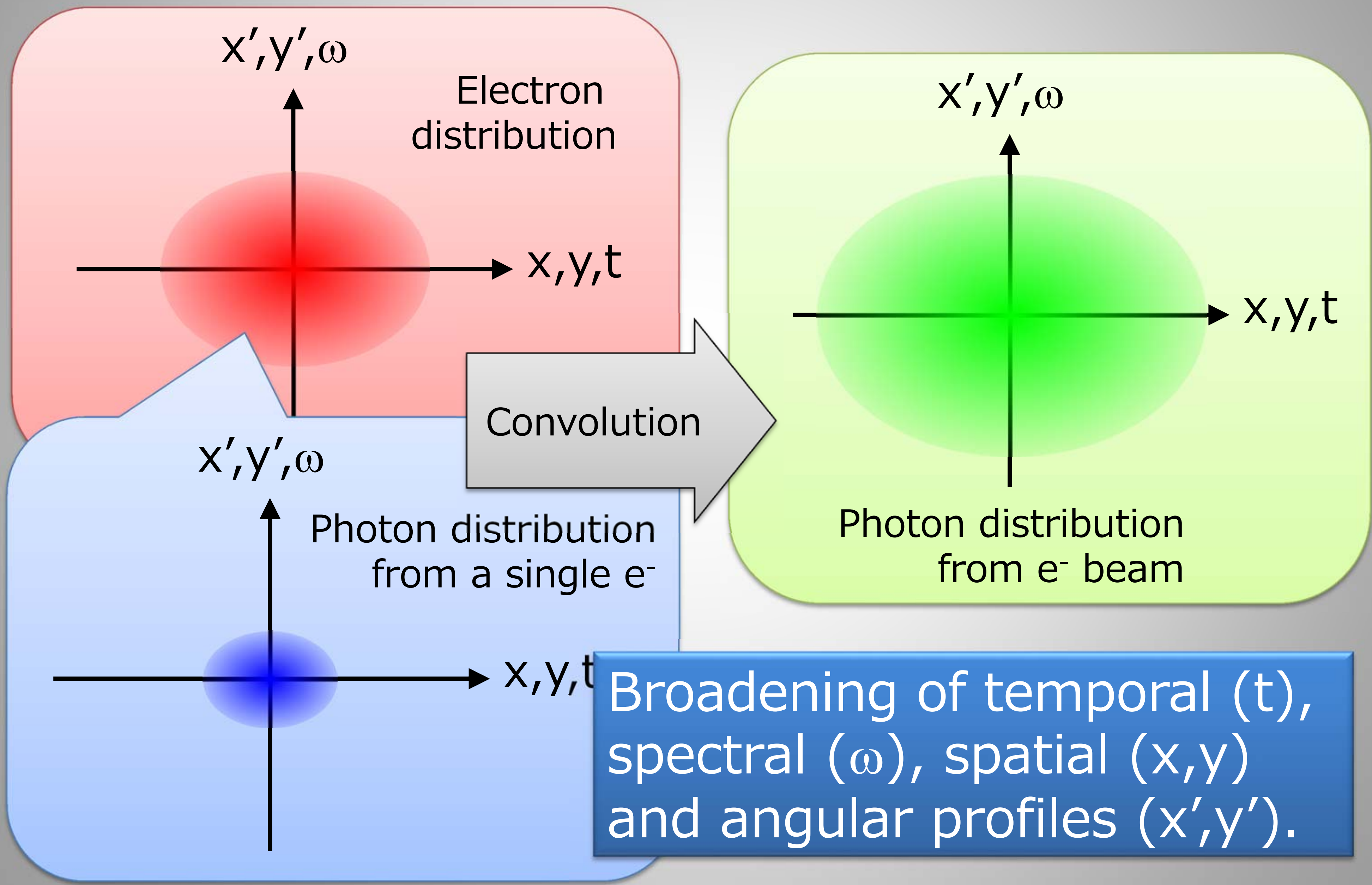
Effective Properties of SR

- Properties of SR emitted from an e^- beam are different from those from a single e^-.
- They are referred to as "effective" properties of SR, as opposed to "natural" properties.

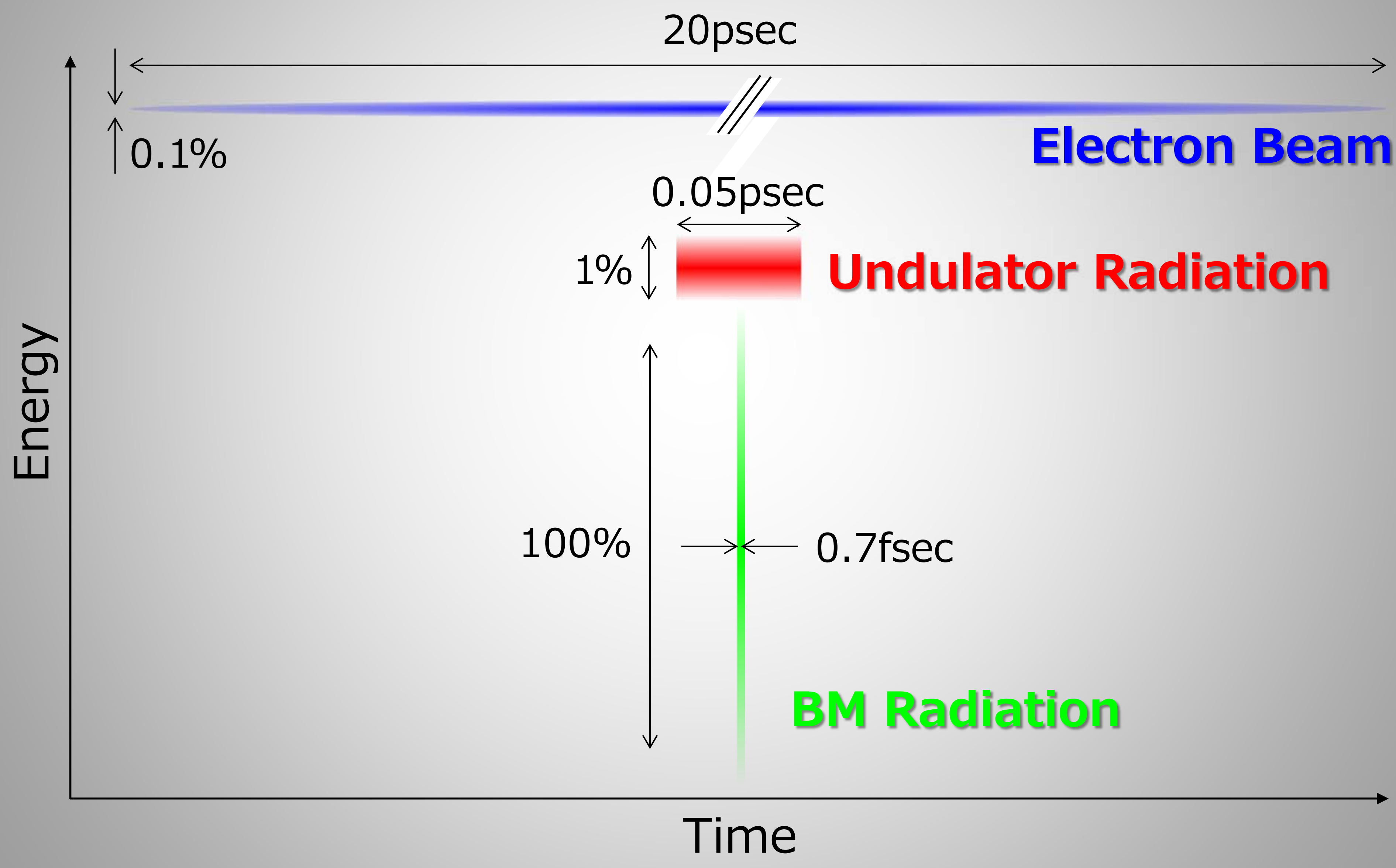


Convolution Between e- and Photon

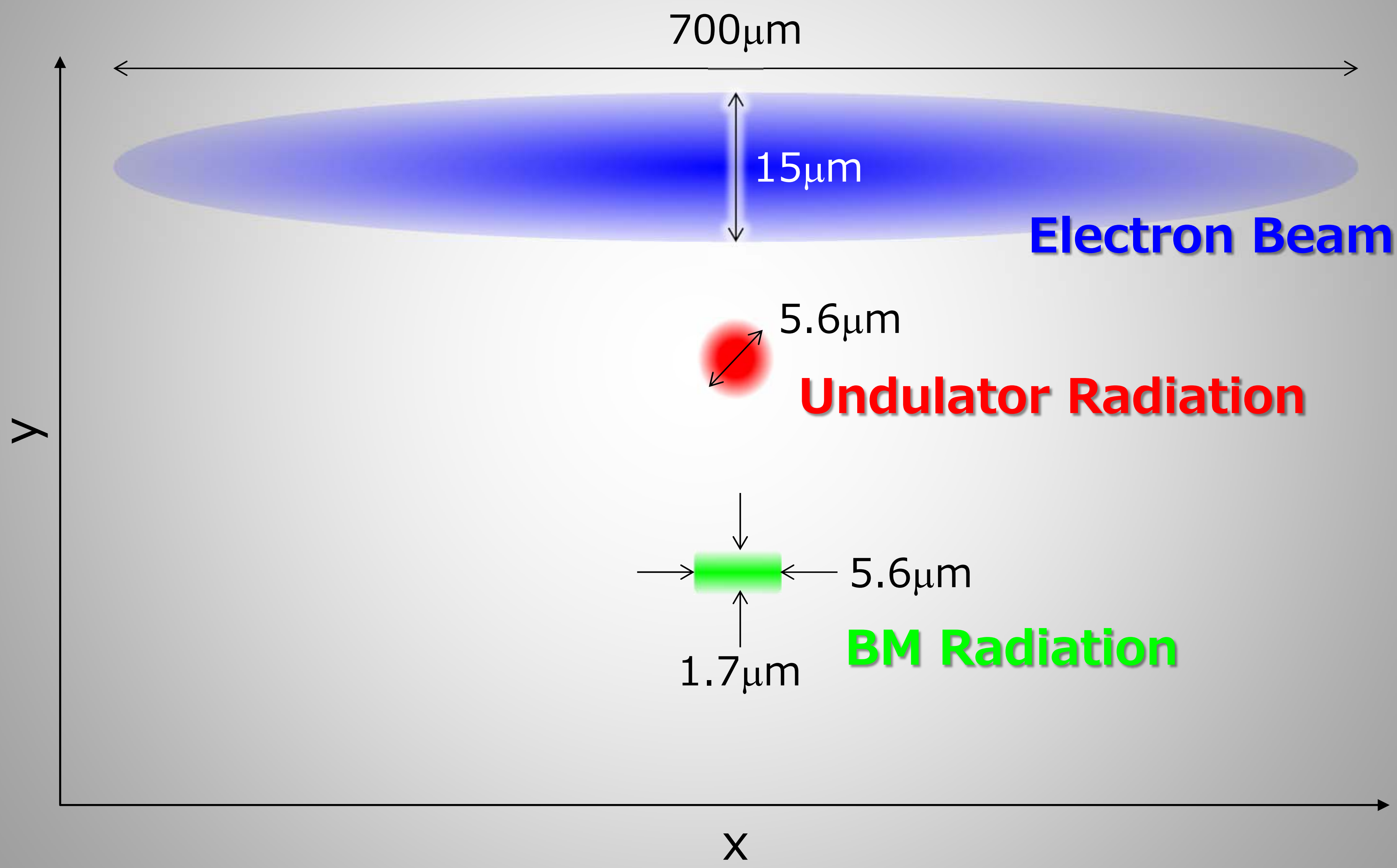


Broadening of temporal (t), spectral (ω), spatial (x,y) and angular profiles (x′,y′).

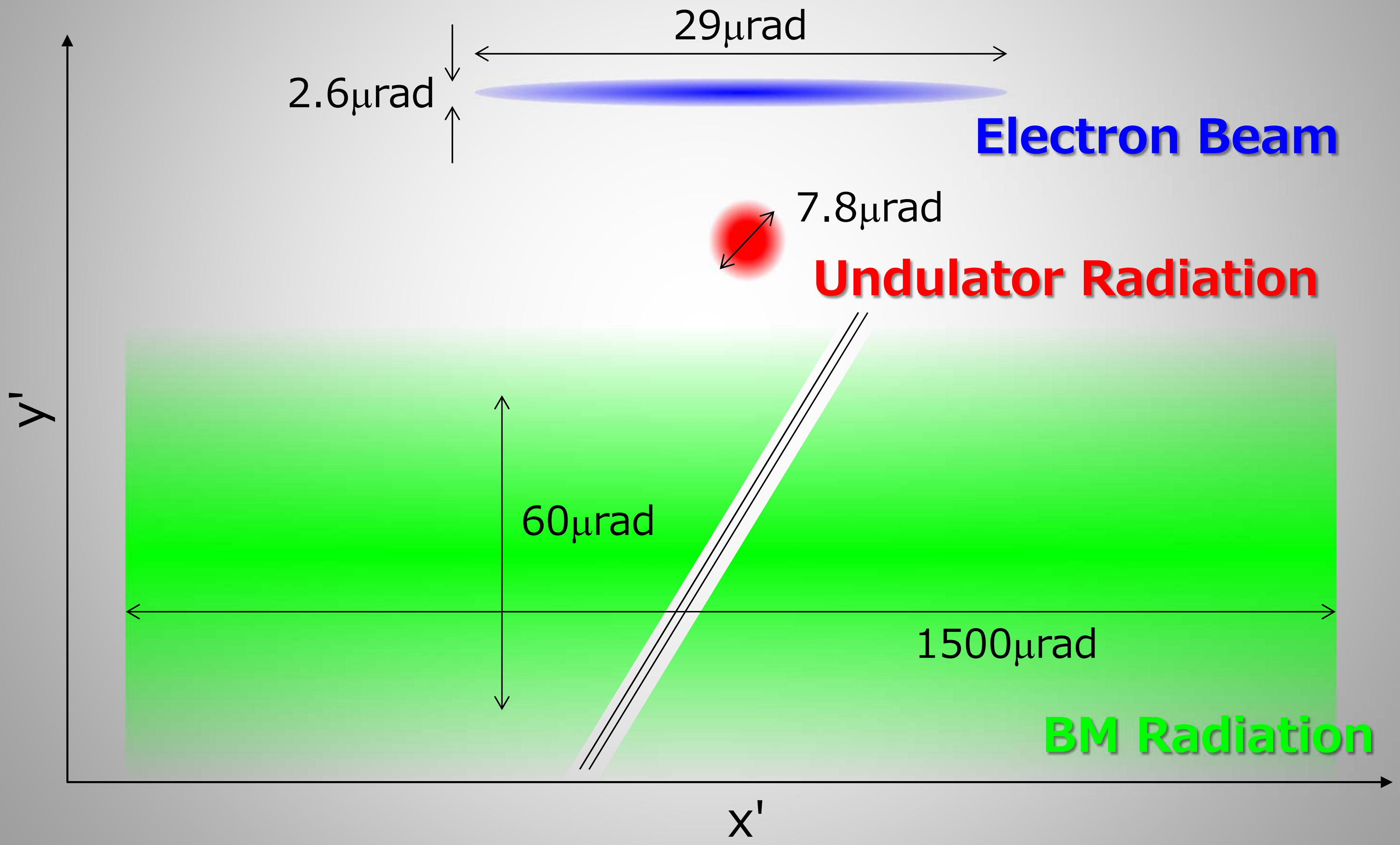
Example in SPring-8: E-t Phase Space



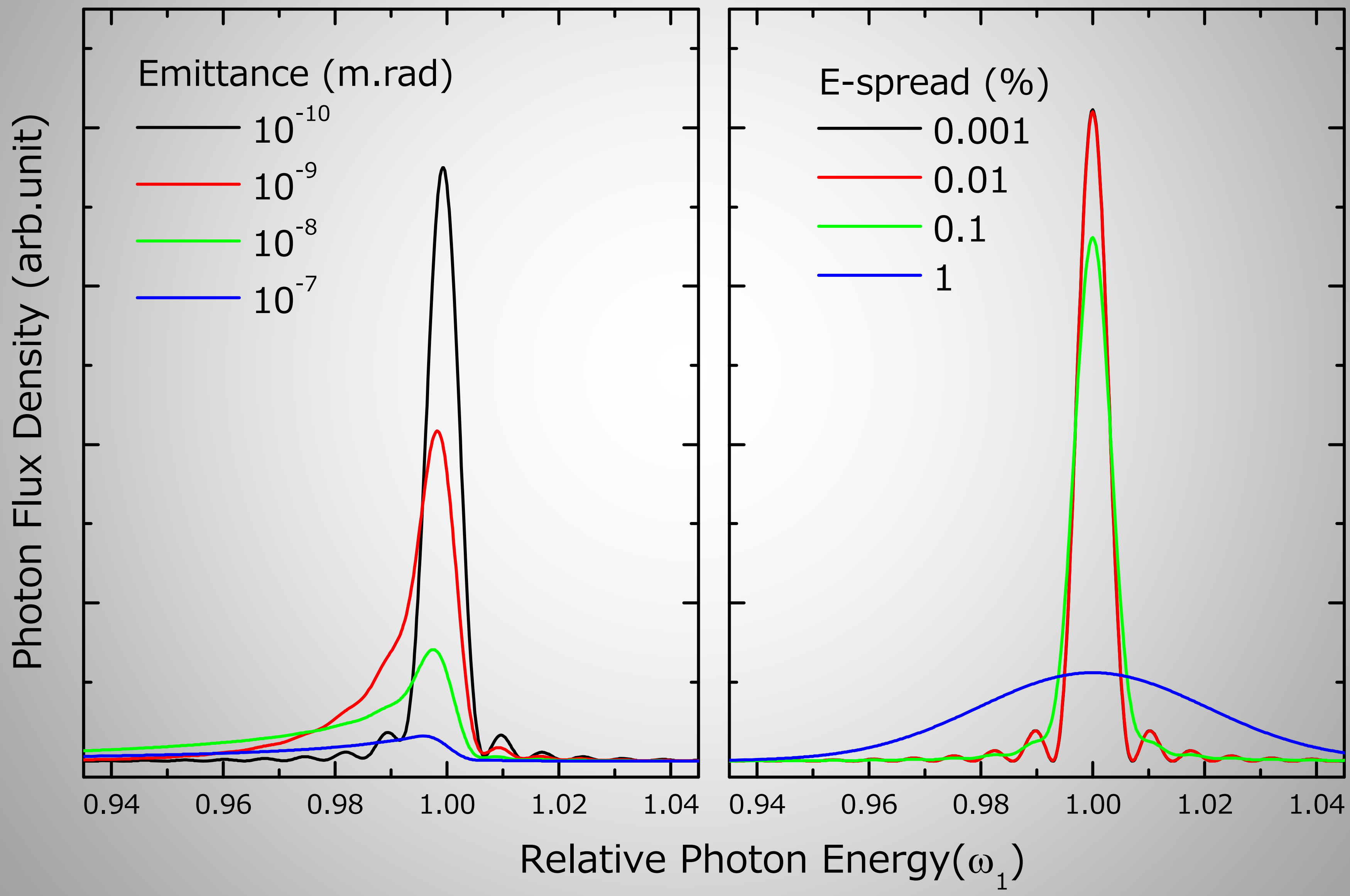
Example in SPring-8: (x,y) Space



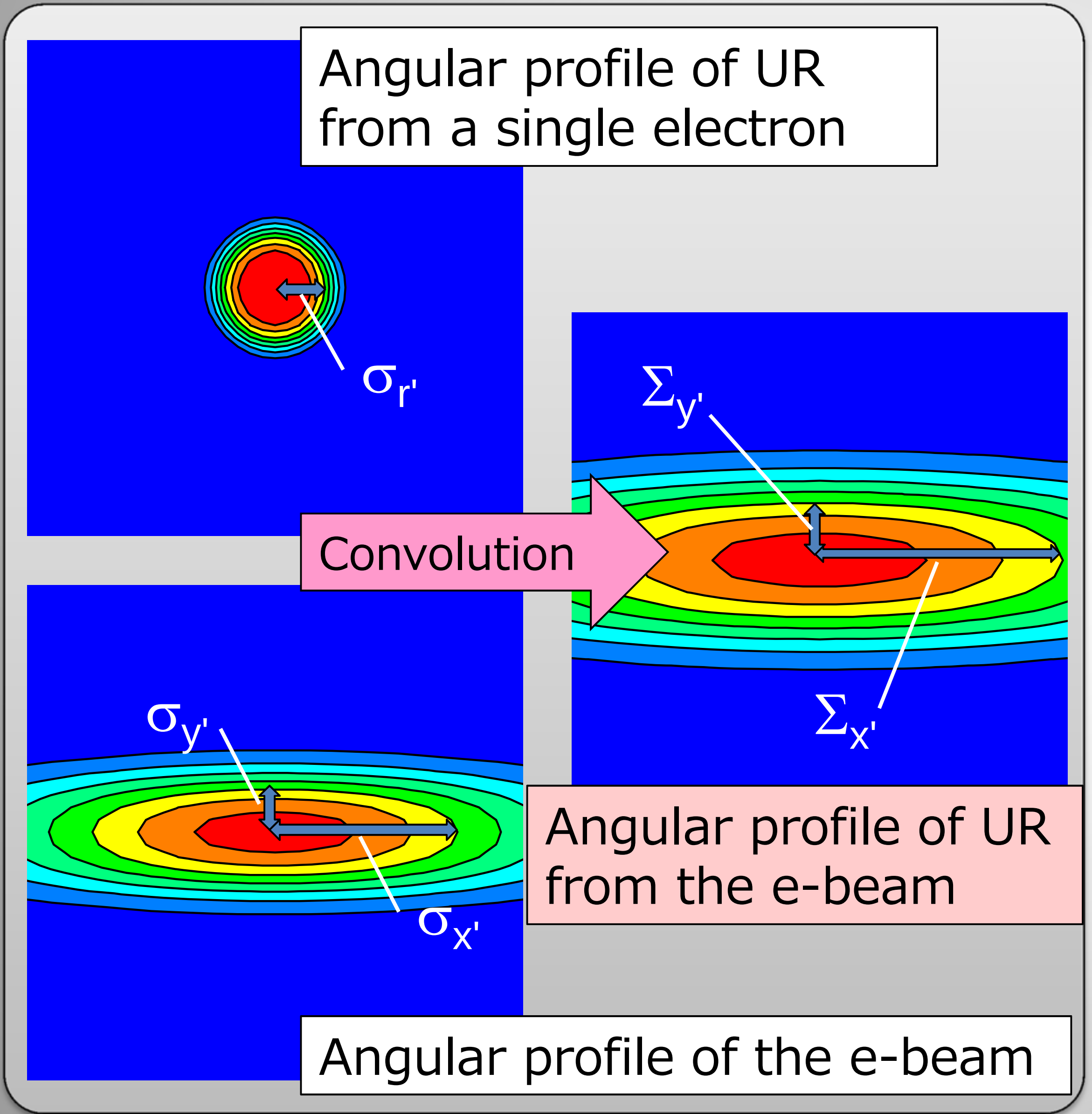
Example in SPring-8: (x′,y′) Space



Spectral Profile (UR)



Angular & Spatial Profile (UR)



✓Gauss approximation
✓Convolution theorem

$$\Sigma_{x',y'} = \sqrt{\sigma_{r'}^2 + \sigma_{x',y'}^2}$$

Effective Angular Div.

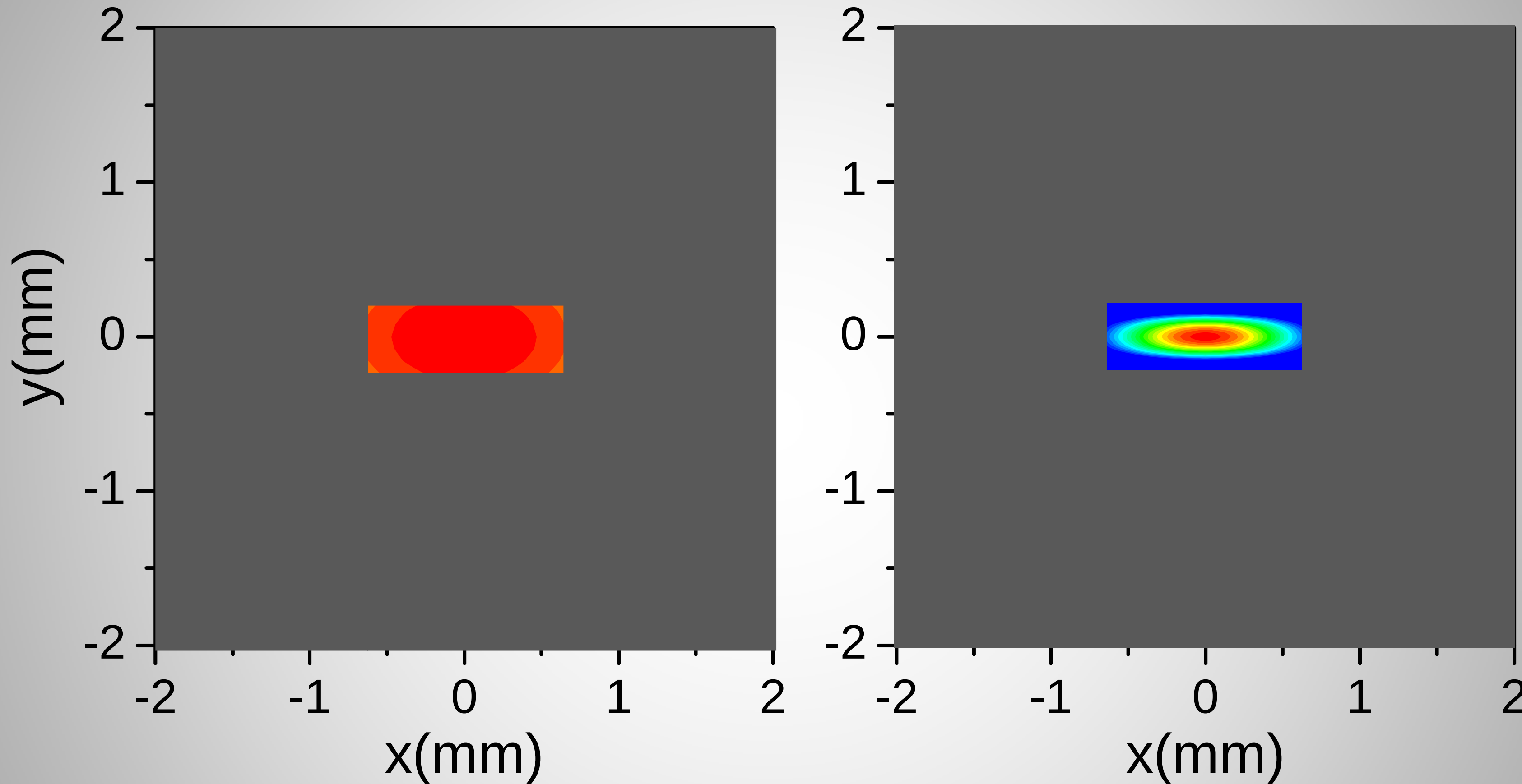
$$\Sigma_{x,y} = \sqrt{\sigma_r^2 + \sigma_{x,y}^2}$$

Effective Source Size

Heat Load on Optical Elements

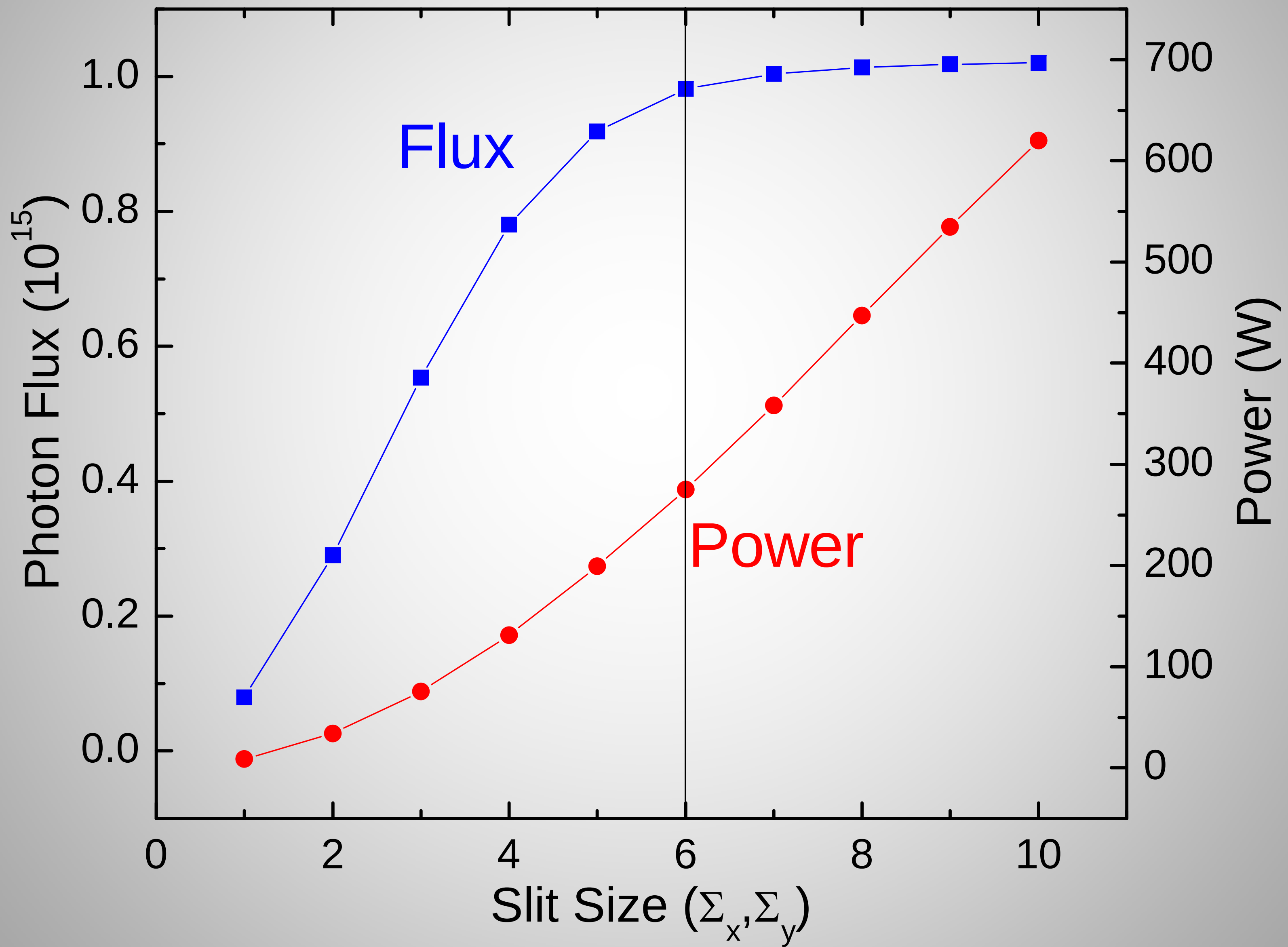
- SR is usually processed by several optical elements before irradiation to the sample, such as the focusing mirror, monochromator.
- These elements can be easily damaged by the heat load of SR.
- In the case of UR, the heat load can be reduced by taking advantage of the difference in the angular profile of the photon flux and radiation power.

Spatial Profile of Power and Flux (UR)



The power profile is much broader than the flux.
Extraction of SR with an appropriate slit is thus effective.

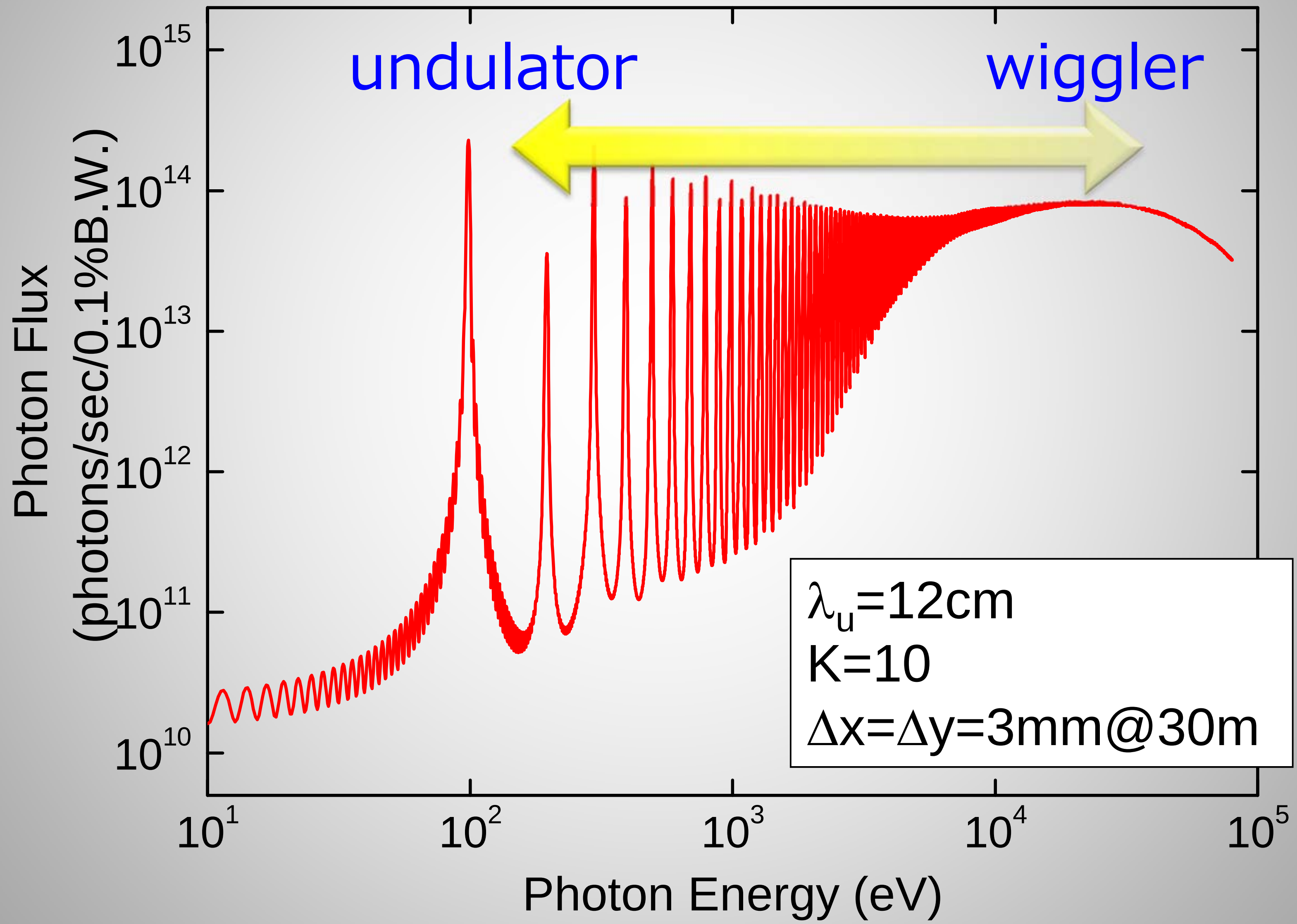
What’s the Optimum Aperture Size?



Wiggler? Undulator? (1)

- Wigglers are identical to undulator from the point of view of magnetic circuit; then what is the distinction?
- It is generally said that the K value distinguishes between the two, however, this is not exactly correct.
- What we should take care is the photon energy region of interest.

Wiggler? Undulator? (2)

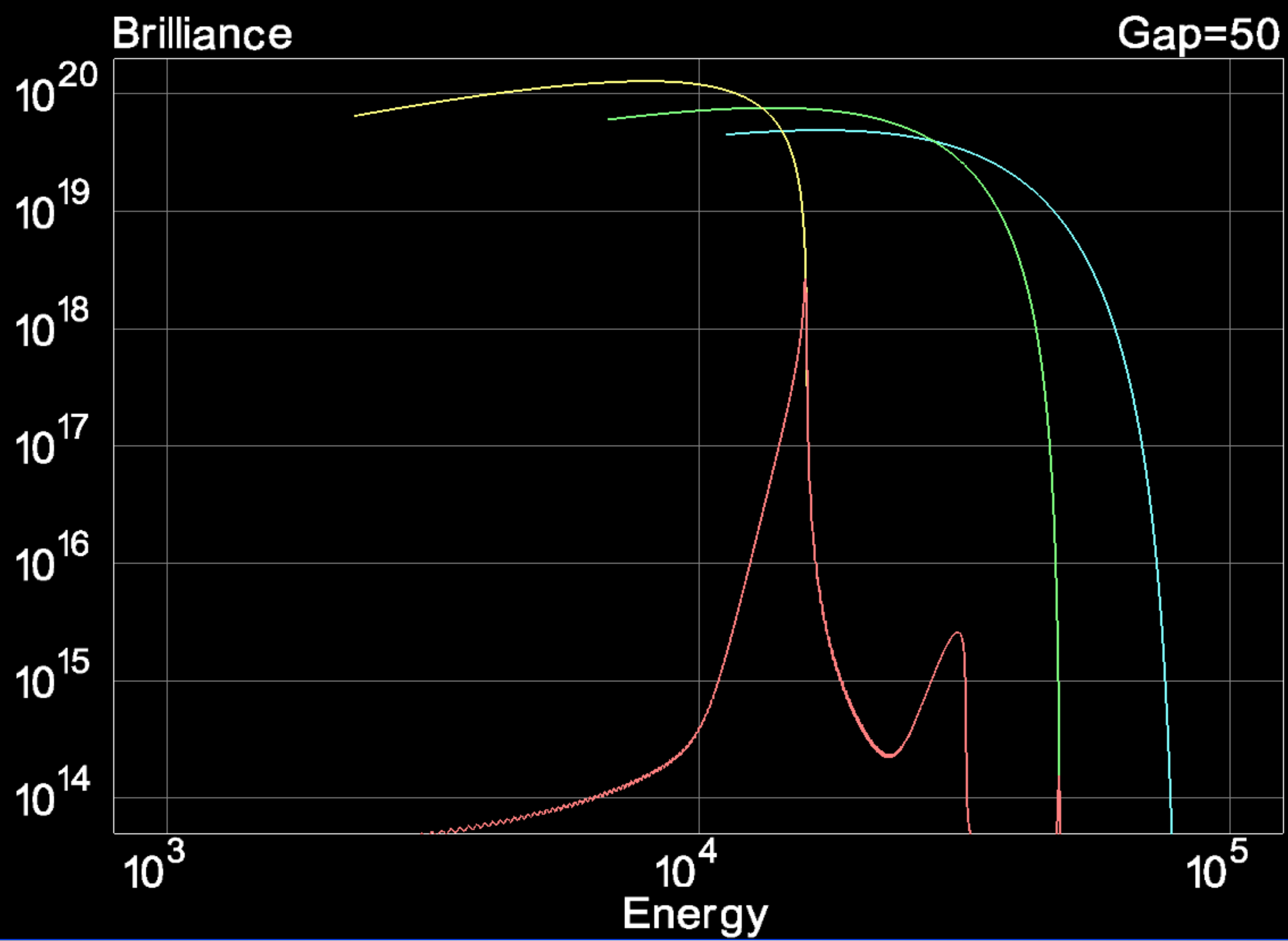


Undulator Radiation Gallery

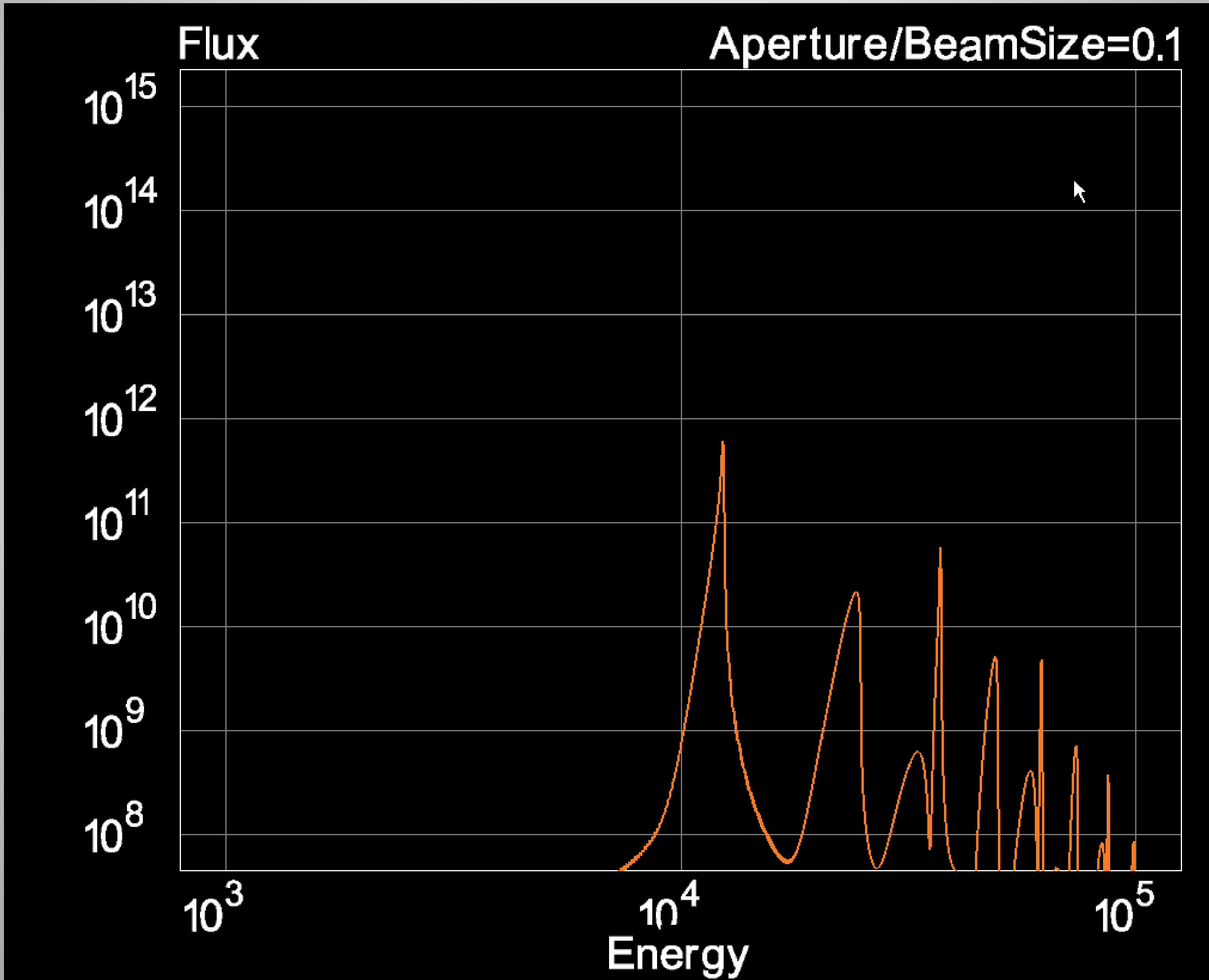
- For quantitative evaluation of SR, a computer code “SPECTRA” is available.
- SPECTRA also offers a function to “visualize” the computation results for further understanding of SR.
  - brilliance curve & spectrum
  - on- and off-peak angular profiles of flux
  - on- and off-axis spectra
  - effects of opening the slit aperture
  - undulator-to-wiggler transition

Brilliance Curve & Spectrum

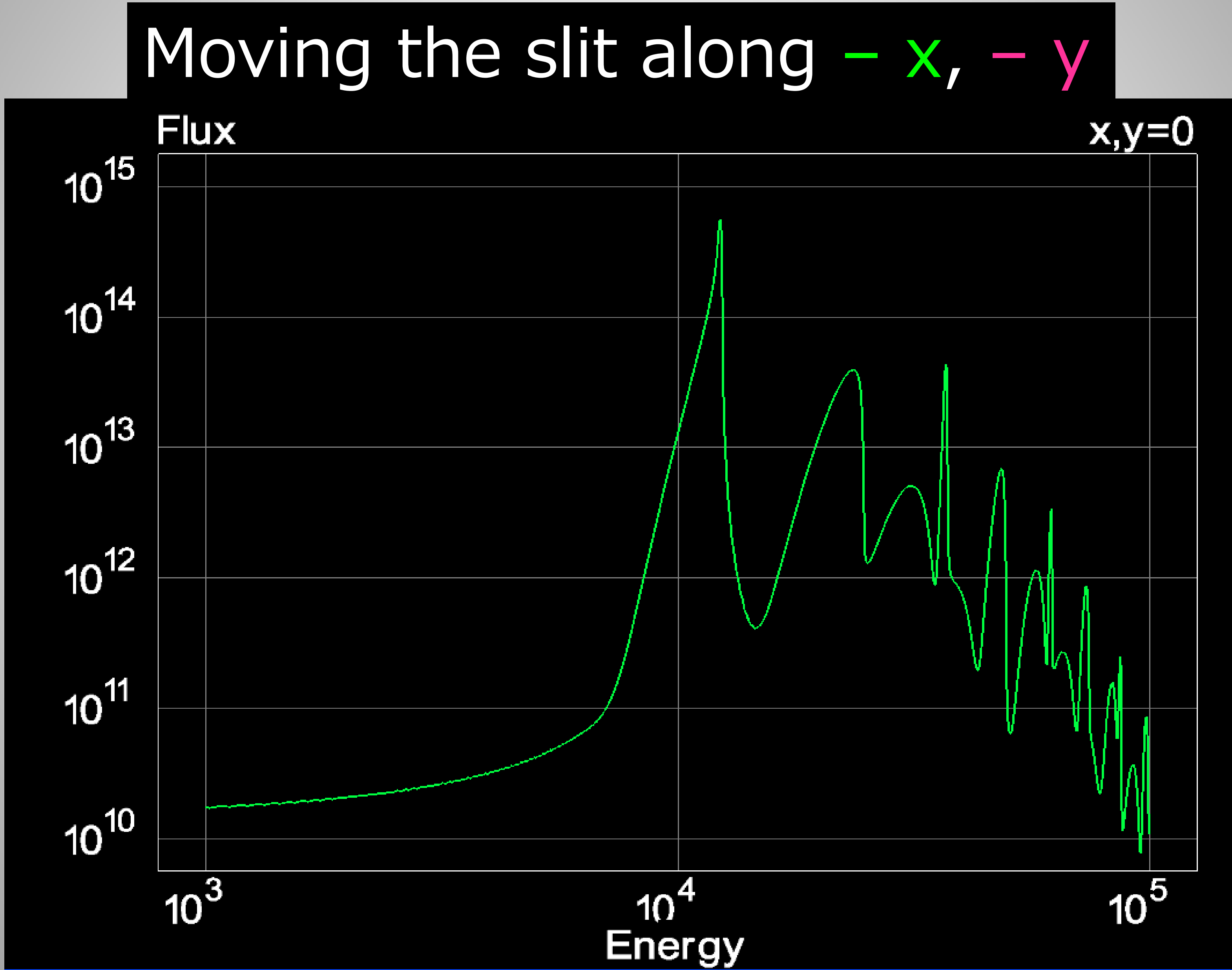
Spectrum –, Peak Brilliance 1st – 3rd – 5th –



Opening the Slit Aperture



Off-Axis Spectrum

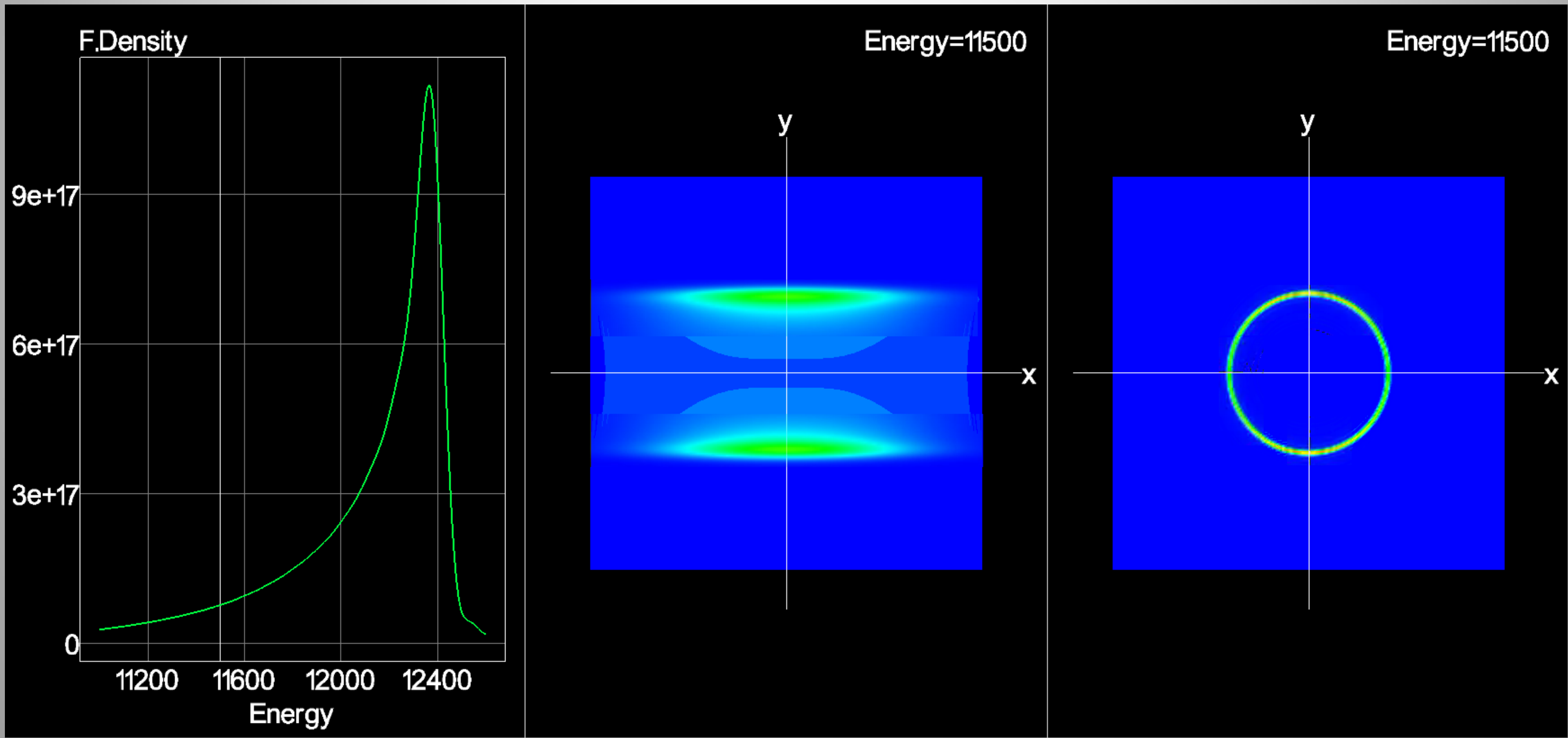


Flux Angular Profile

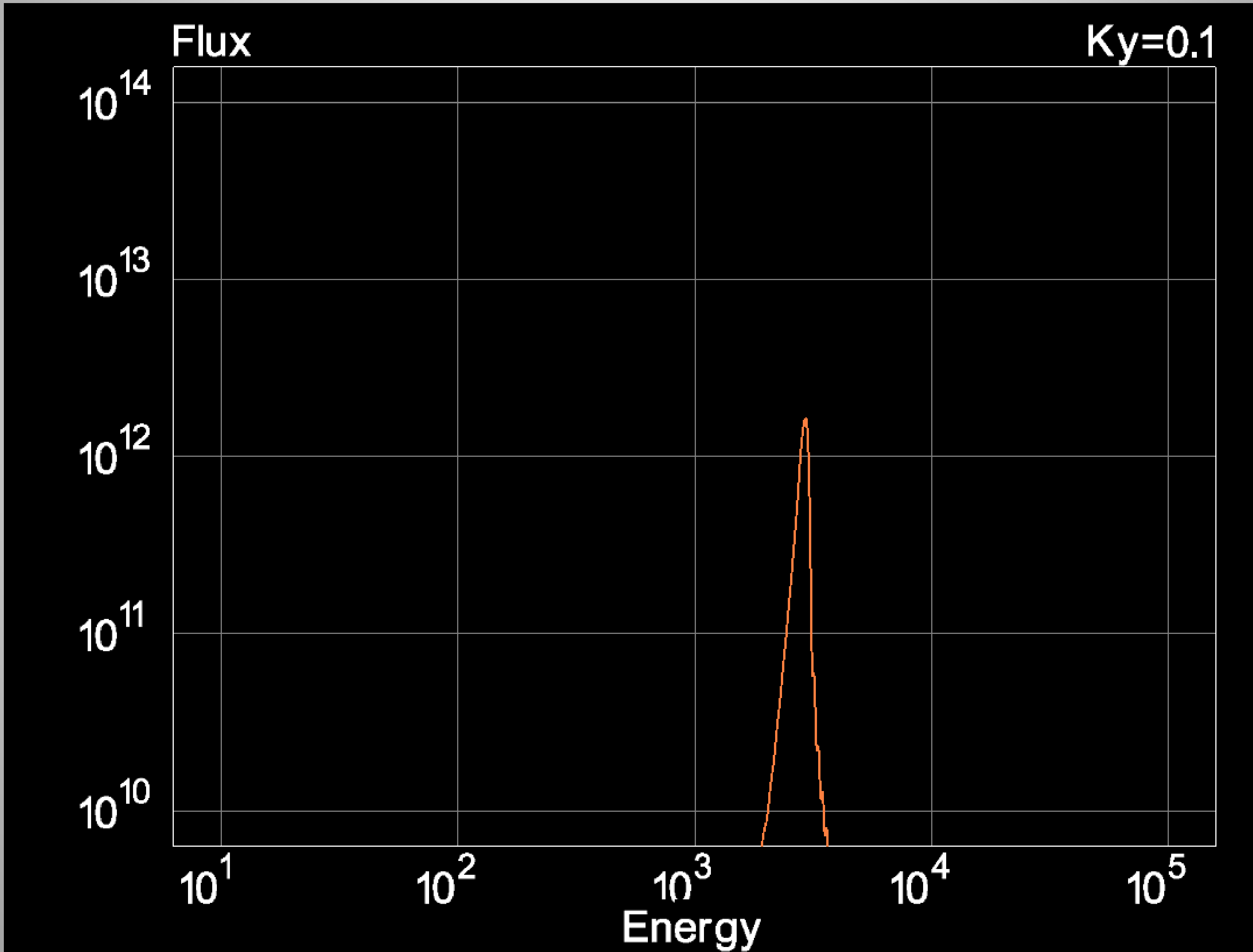
On-Axis Spectrum

Angular Profile
(Finite Emittance)

Angular Profile
(Zero Emittance)



Undulator-to-Wiggler Transition



Other Important Topics

- Polarization related issues
- Quantitative descriptions of SR
- Light sources for CPR & helicity switching
  - helical undulator & elliptic wiggler
  - chicanes & choppers, kicker magnets
- Effects on the electron beam
  - natural focusing
  - beam-axis fluctuation due to COD variation
- Undulators for SASE-based X-ray FEL

Announcement

- Those who are going to join the lecture on "SPECTRA" scheduled on 16th (Wed.) evening are recommended to bring your own PC.
- If you do not have any PC, please consult the office.
- Before joining the lecture, download the software from the URL below and finish installation.

http://radiant.harima.riken.go.jp/spectra